



Search for Neutral MSSM Higgs Bosons in the Di-Tau Decay Channel at CDF

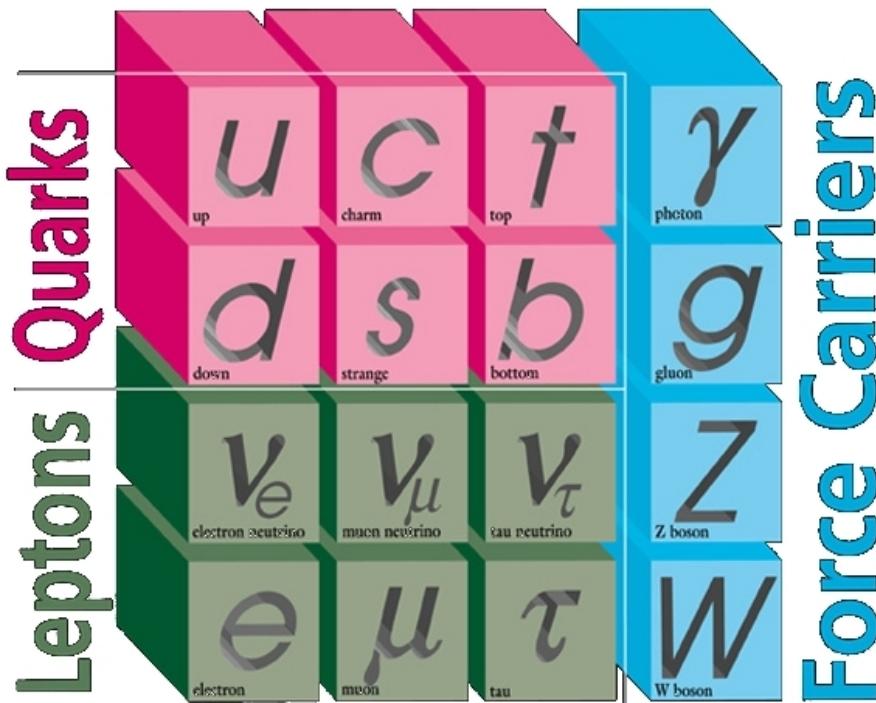
Anton Anastassov
(Rutgers University)

Particle Physics Seminar
BNL - February 15, 2007

Contact: aa@fnal.gov

- Searching for the Higgs boson at the Tevatron
- SM vs MSSM Higgs
- The tau lepton as a tool
- Taus at CDF
- Preliminary results with $\mathcal{L}_{\text{int}} = 1 \text{ fb}^{-1}$
- Summary and outlook

The Standard Model – what is missing?

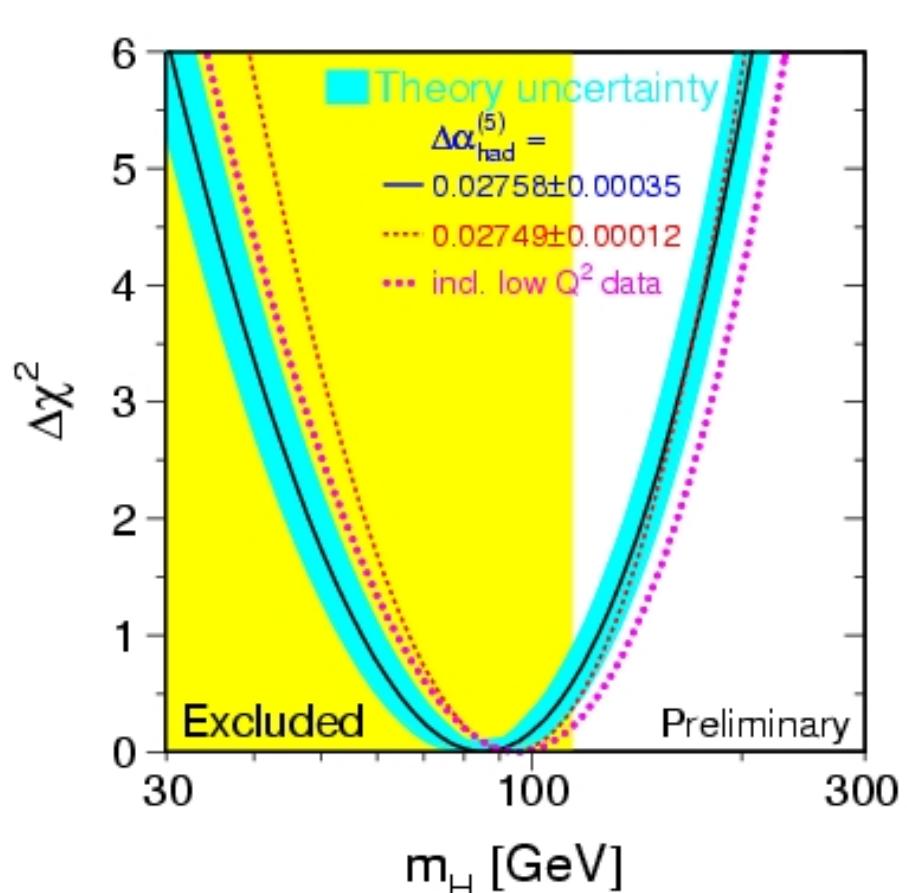


- Good description of particles and their interactions
- Extensively tested
- But... need explanation of W,Z, fermion masses
- Suggested solution: the Higgs mechanism
 - Predicts the existence of a new massive neutral particle – the Higgs boson

Understanding the origin of masses is one of the priorities at the Tevatron and LHC.

Where is the SM Higgs Boson?

- Theory does not predict the mass of the Higgs:
 - LEP limit: $m_H > 114$ GeV @ 95% CL
 - Indirect limit from EW data:

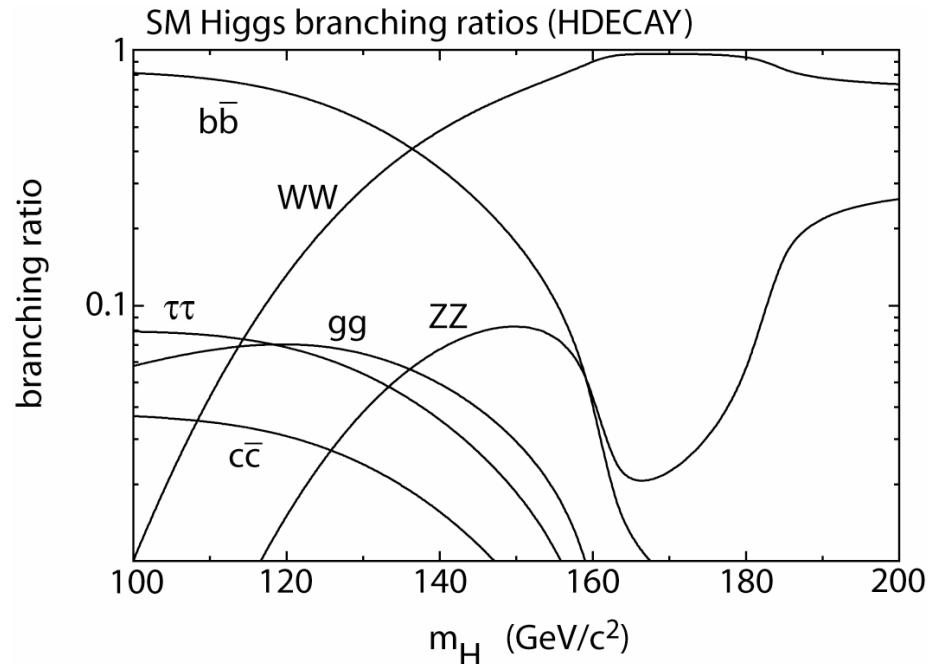
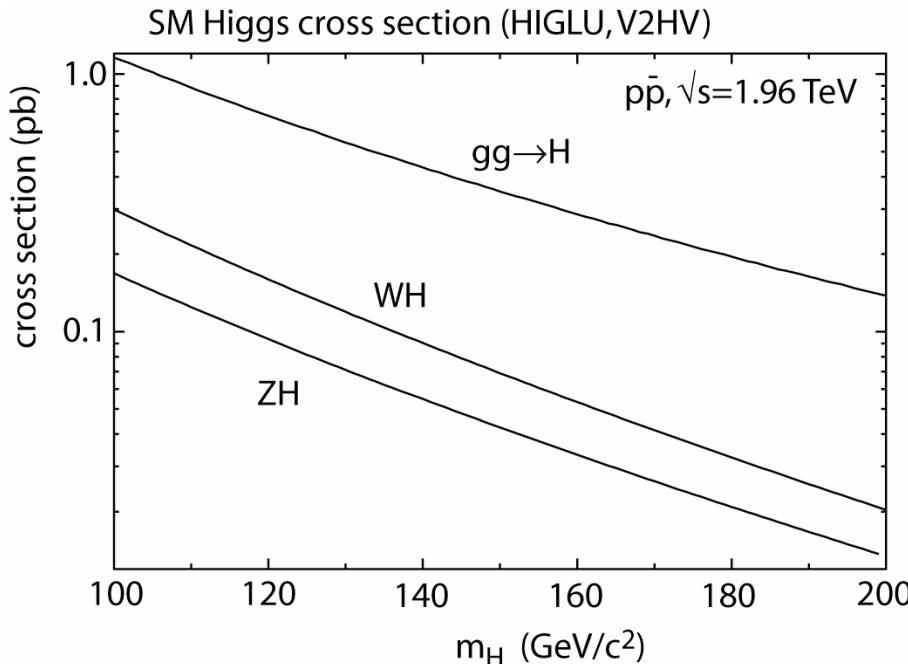


- Preferred value: $m_H = 85^{+39}_{-28}$ GeV
- $m_H < 166$ GeV @ 95% CL (one-sided)
- $m_H < 199$ GeV @ 95% CL (including LEP limit from the direct search)

Results as of Summer 2006!

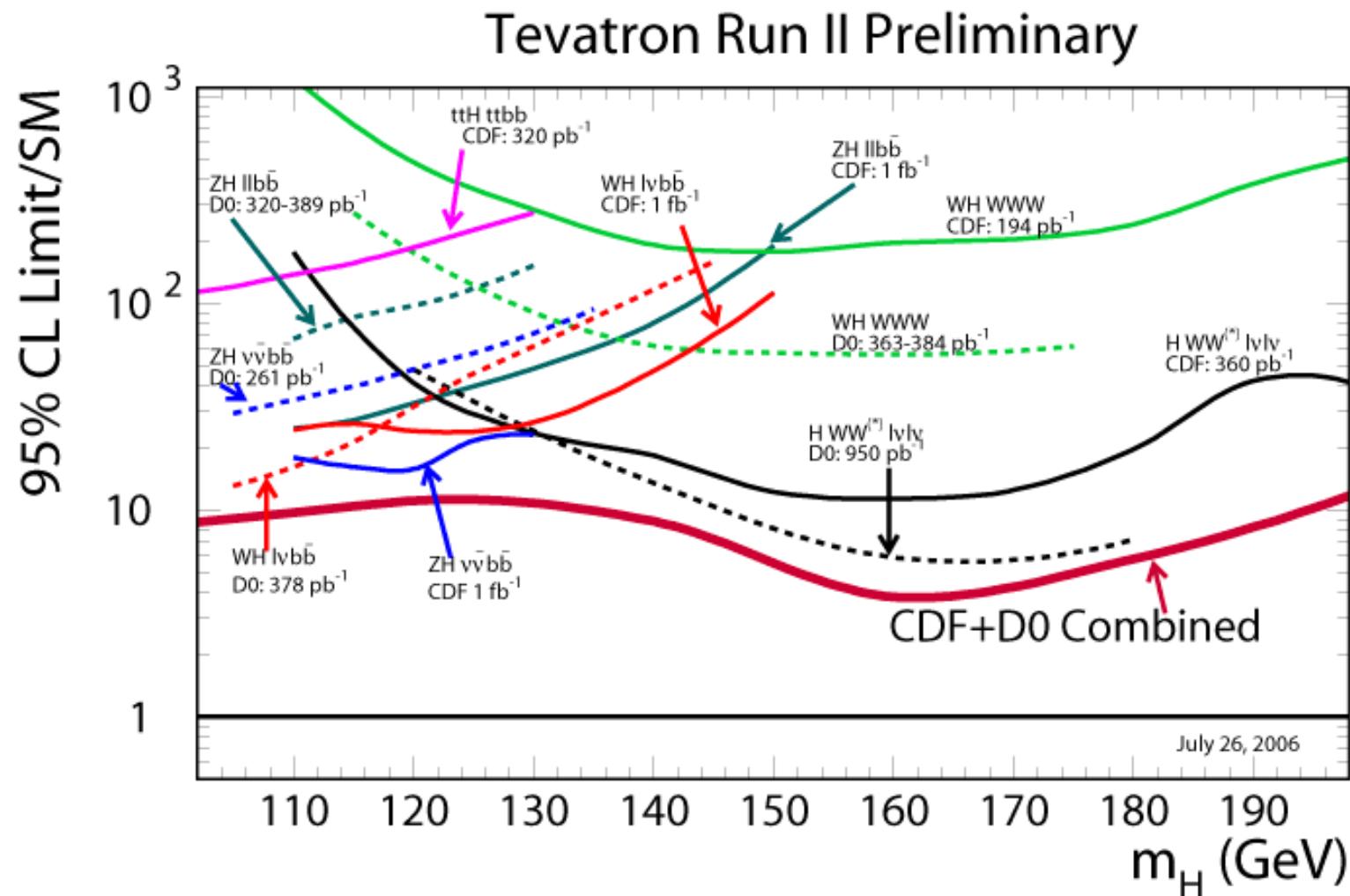
- Relatively light Higgs is favored
- The new CDF measurement of m_W pulls the estimate of m_H even lower

How to search for SM Higgs at the Tevatron



- Most promising for lower mass-range: associated production with W/Z; $W \rightarrow l\nu$, $Z \rightarrow ll, \nu\nu$; $H \rightarrow b\bar{b}$
- Higher masses: $gg \rightarrow H$; $H \rightarrow WW^*$

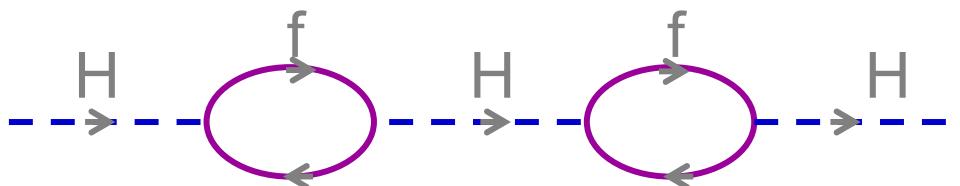
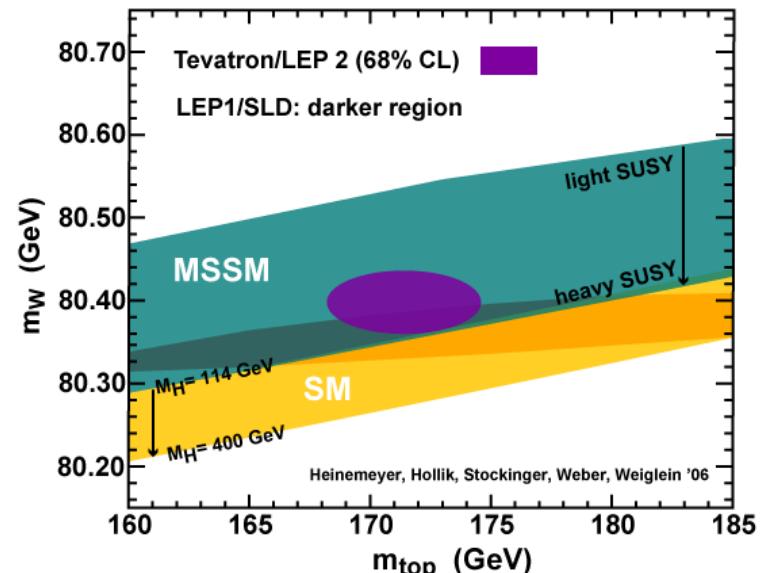
Where we are at the end of 2006



Getting closer to the SM predictions... but not there yet

Why Search for Higgs Beyond SM?

- Possible enhancement in SUSY Higgs production
- Recent results seem to prefer SUSY?
- Difficulties in the SM Higgs sector:
 - Large corrections to Higgs mass: $\delta m^2 \sim \Lambda^2$; $\Lambda = M_P$?



- Special tuning needed (“accidental” cancellation?)

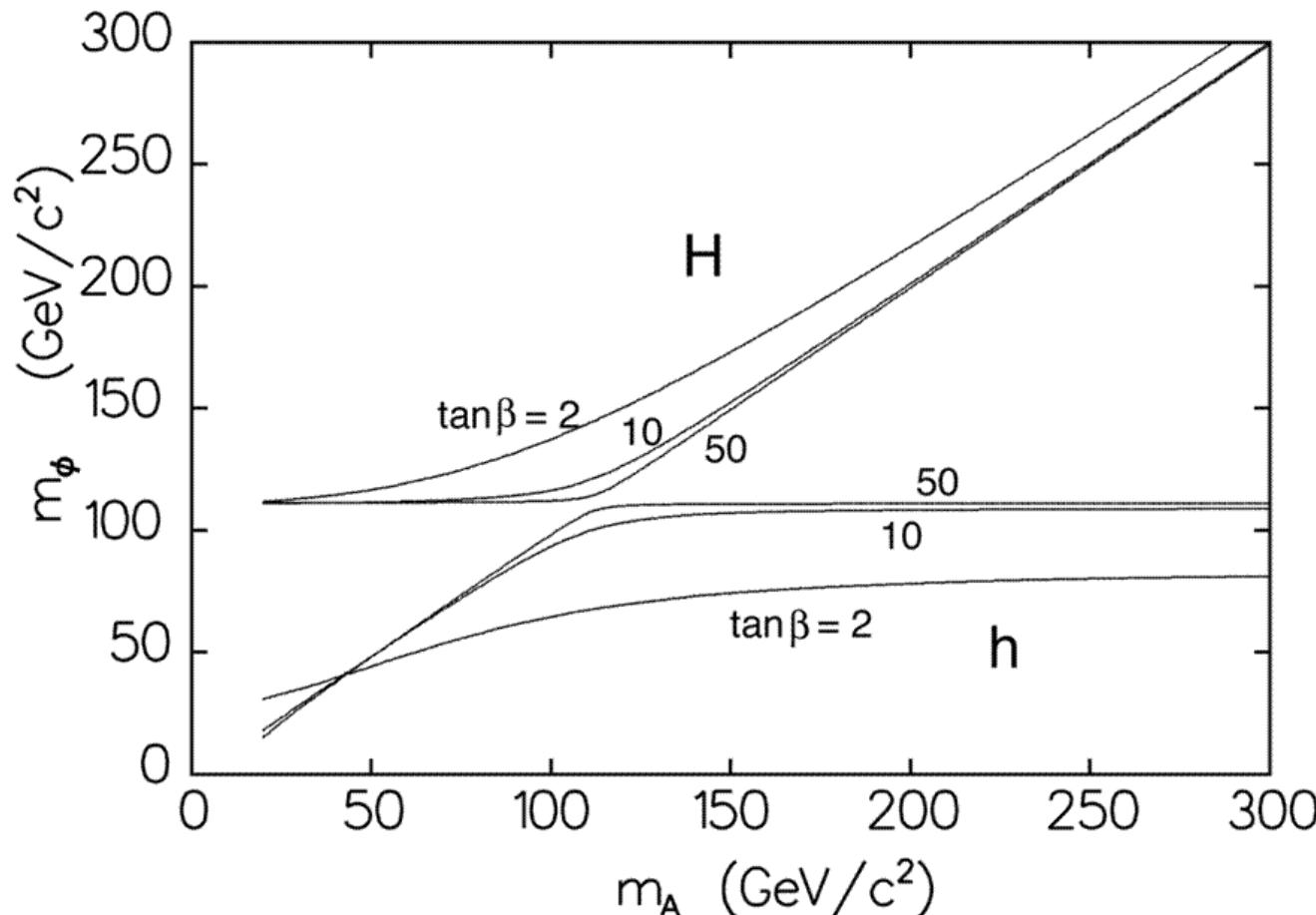
SUSY provides an elegant solution

- New symmetry: fermions (bosons) have bosonic (fermionic) partners
 - Doubles the number of particles
- Unification of forces at one scale
- Dark matter candidate (LSP)
- ...
- **As far as Higgs is concerned:**

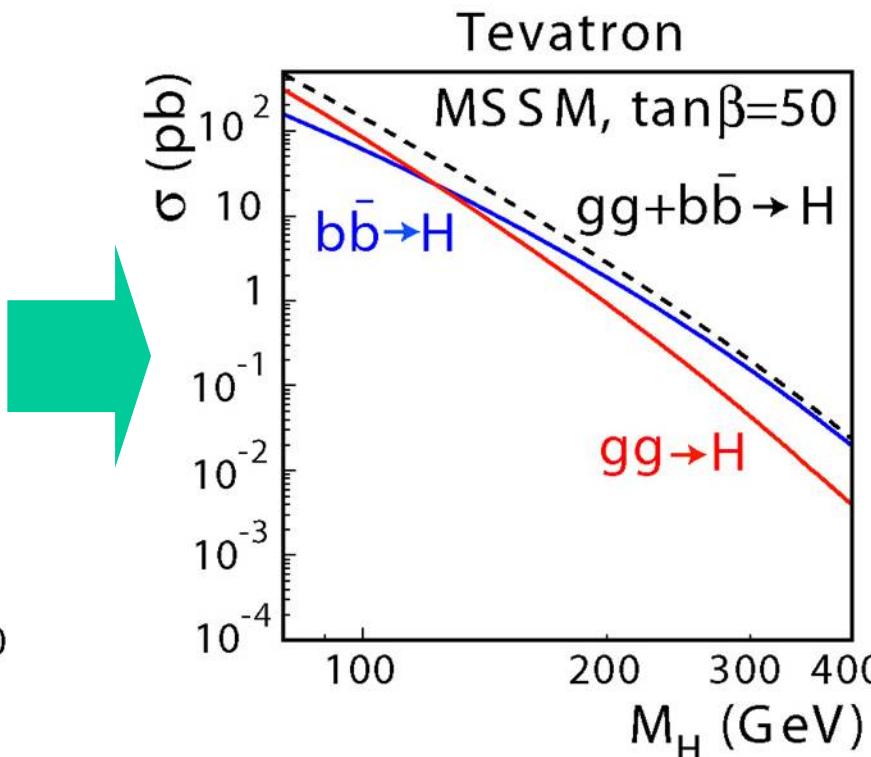
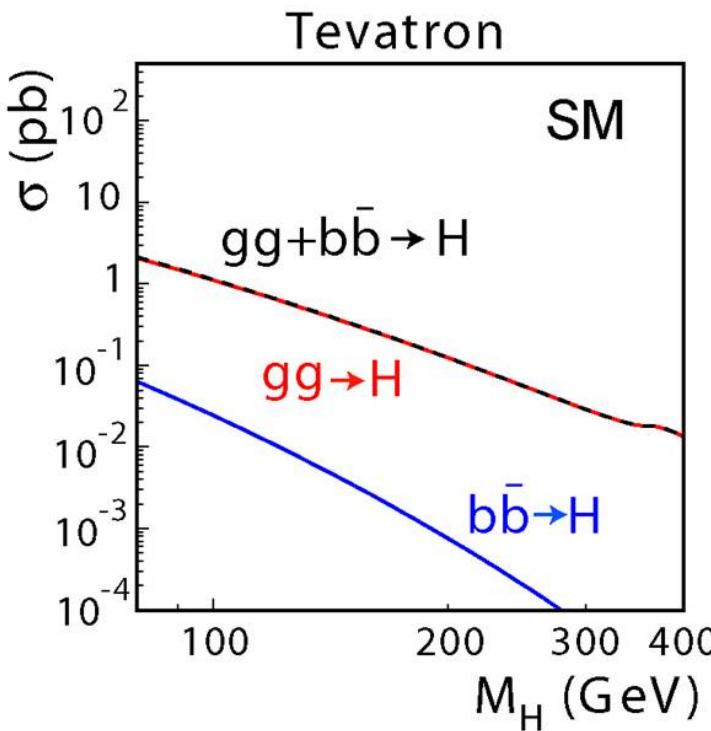
Divergent terms in m_H calculations canceled by the spartners – the “accidental” cancellation has solid underlying foundation

Minimal Supersymmetric Standard Model (MSSM): SUSY extension of SM with minimal particle content

- Higgs sector in the MSSM:
 - Requires two Higgs field doublets
 - Five physical states: $H, h, A; H^\pm$
 - Lightest Higgs (h) mass close to EW scale
 - At tree level defined by m_A and $\tan\beta = v_u/v_d$
 - A couplings to b, τ enhanced by $\sim \tan\beta$!



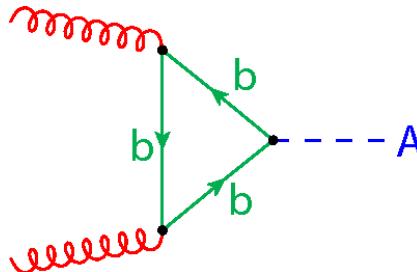
- For large $\tan\beta$ h or H are almost mass degenerate with A , similar couplings
- The other one is SM-like, low-mass ($m < 135$ GeV)



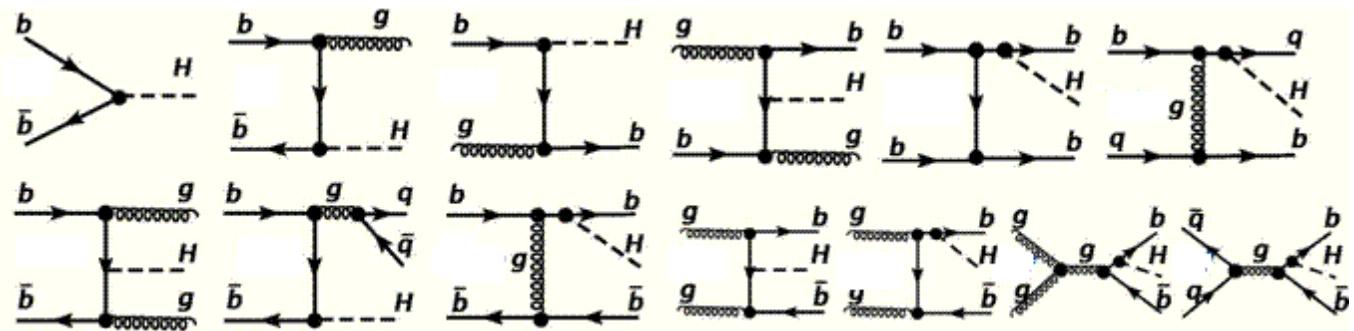
Higgs decays:

- bb ($\sim 90\%$)
- $\tau\tau$ ($\sim 10\%$)

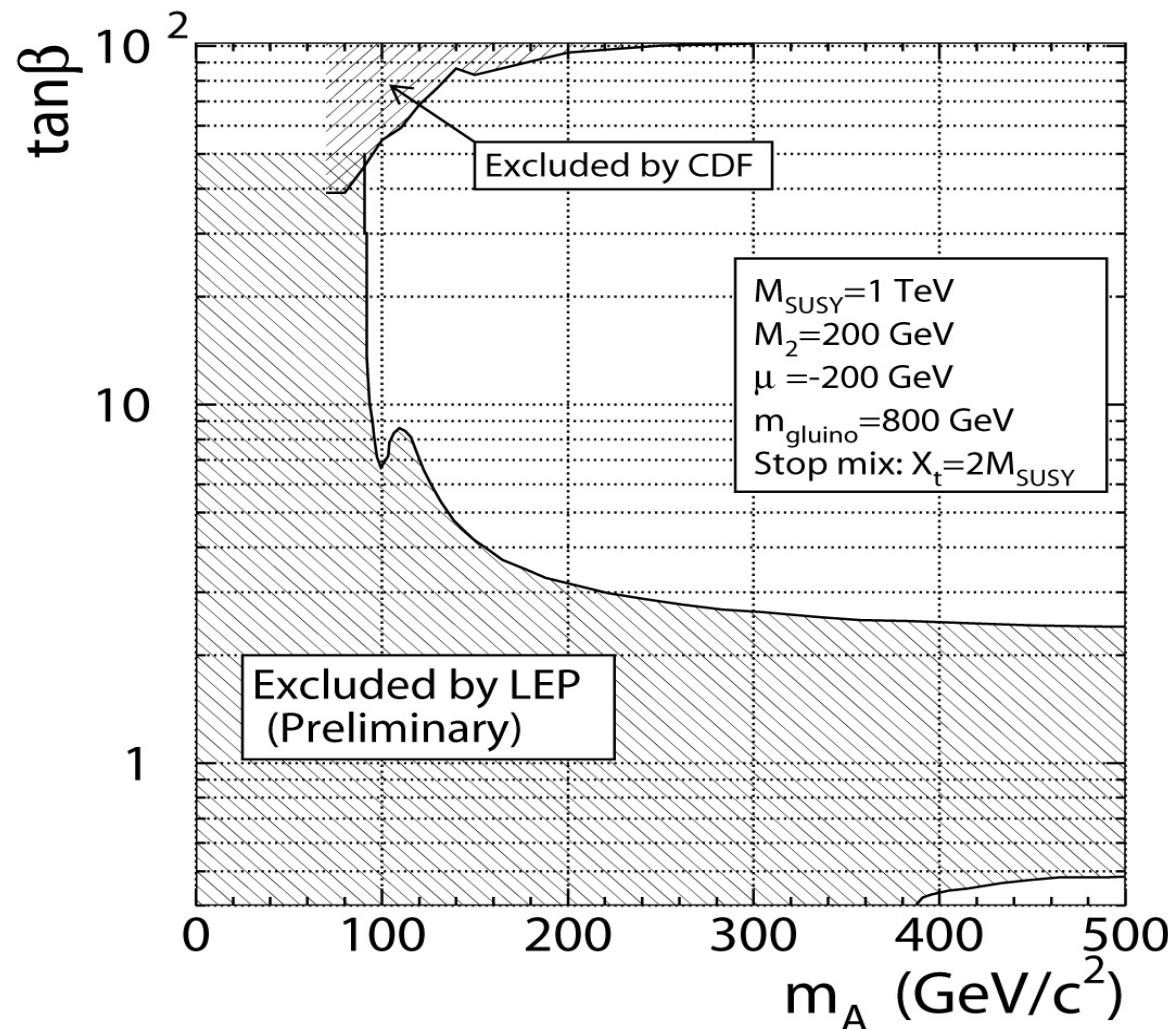
$gg \rightarrow H$



$b\bar{b} \rightarrow H$



Neutral MSSM Higgs Limits: before Run II



Large regions of parameter space remained unexcluded.

The $A+b(b) \rightarrow bbb(b)$ channel was considered the most promising search mode
(Tevatron Higgs Working Group Report).

- BR($A \rightarrow bb$) is larger, *however*, the strengths of a search channel are influenced by other factors as well
- Due to large QCD backgrounds, the searches for $A \rightarrow bb$ are restricted to production mechanisms with an associated high- p_T b-jet: $A+b(b)$
 - Can not take advantage of $gg \rightarrow A$
 - Explore only a subset of the $bb \rightarrow A$
- All production mechanisms can be explored with the $A \rightarrow \tau\tau$ channel
 - **Largely compensates the BR disadvantage!**

Case for MSSM Higgs $\rightarrow\tau\tau$ Search

- MSSM Higgs production and decays can be significantly affected by radiative corrections *
- The $b\bar{b}$ channel is more sensitive to these corrections (and therefore to the SUSY scenario), while the $\tau\tau$ channel is more robust

$$\sigma(b\bar{b} \rightarrow A) \times BR(A \rightarrow b\bar{b}) \cong \sigma(b\bar{b} \rightarrow A)_{SM} \times \frac{\tan \beta^2}{(1 + \Delta_b)^2} \times \frac{9}{(1 + \Delta_b)^2 + 9}$$

$$\sigma(b\bar{b}, gg \rightarrow A) \times BR(A \rightarrow \tau\tau) \cong \sigma(b\bar{b}, gg \rightarrow A)_{SM} \times \frac{\tan \beta^2}{(1 + \Delta_b)^2 + 9}$$

Δ_b is a function of SUSY parameters

A $\rightarrow\tau\tau$ is important for MSSM Higgs searches:

- adds an extra detection mode
- sensitive to additional production mechanisms
- expands the probed SUSY parameter space

* M. Carena, S. Heinemeyer, G. Weiglein, and C.E.M. Wagner, Eur.Phys.J. C45 (2006) 797-814

The Tools of the Search: Tau Leptons

- Third generation; charged: $Q^\tau = \pm 1$
- Heavy: $m_\tau = 1.777$ GeV
- Decays :
 - Leptonic: $\tau \rightarrow e \bar{\nu}_e \nu_\tau, \tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau$ ($\sim 36\%$)
 - Hadronic: $\tau \rightarrow \pi \nu_\tau, \tau \rightarrow \pi \pi^0 \nu_\tau, \tau \rightarrow \pi \pi \pi \nu_\tau, \tau \rightarrow \pi \pi^0 \pi^0 \nu_\tau \dots$ ($\sim 64\%$)
- rich decay spectrum!

Taking advantage of $A \rightarrow \tau\tau$ comes at a price:
reconstruction of many and sometimes complex
final states!

Throughout this talk we use $\tau_e, \tau_\mu, \tau_{had}$ as shorthand notations for $\tau \rightarrow e \bar{\nu}_e \nu_\tau, \tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau$, and $\tau \rightarrow hadrons \nu_\tau$, respectively.

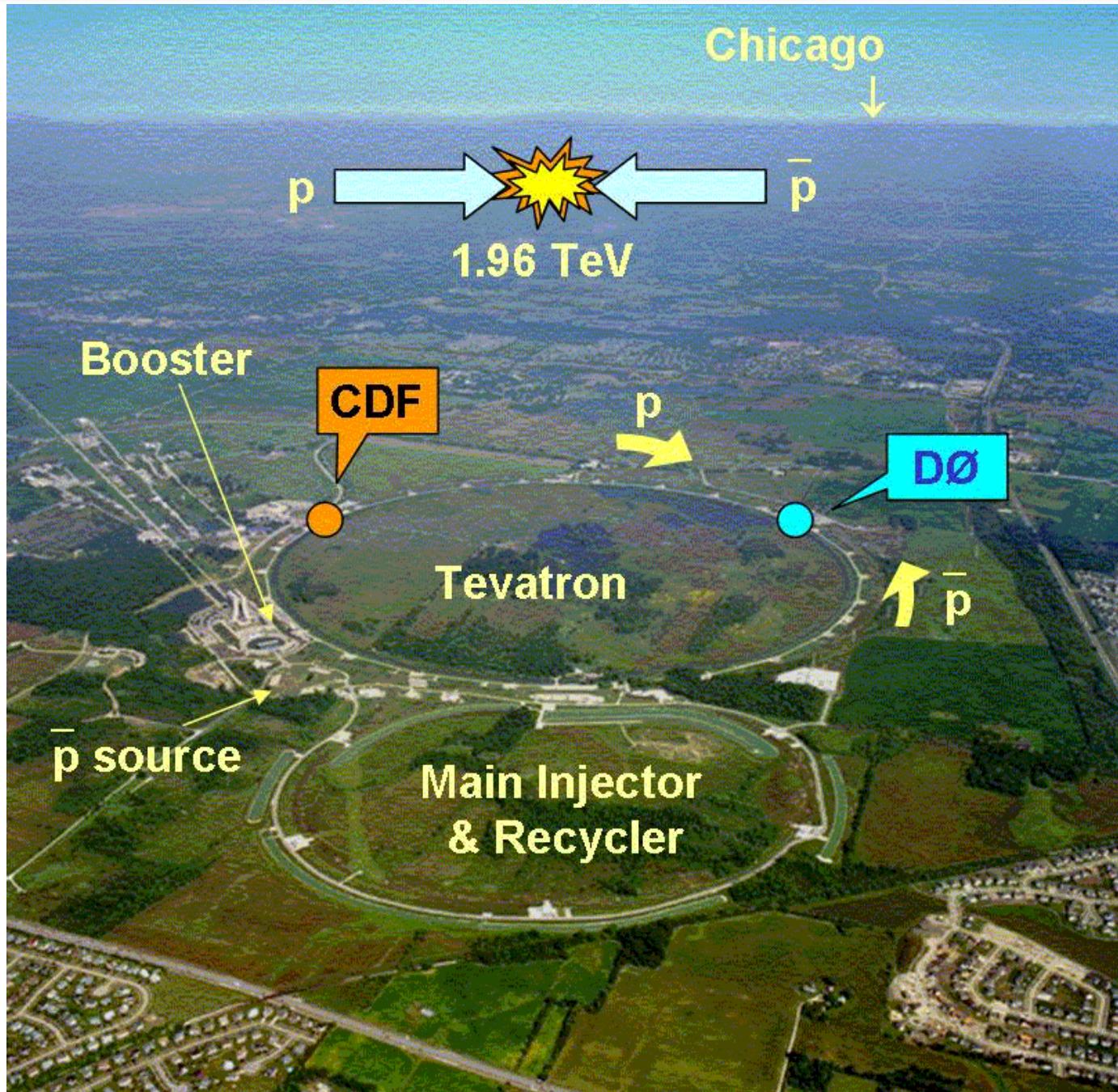
- Determined by the decay modes of the taus

Used in this search {

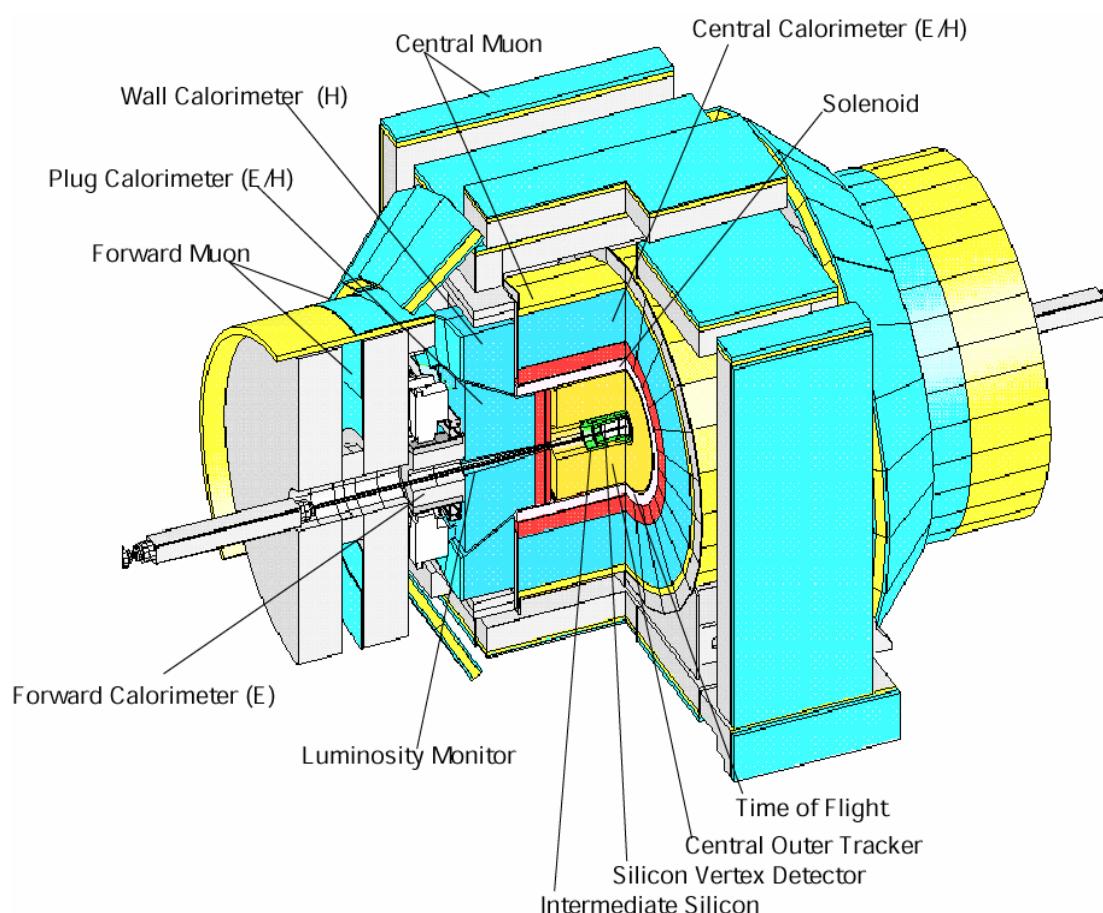
Mode	Fraction (%)	Comment
$\tau_e \tau_e$	3	large $Z/\gamma^* \rightarrow ee$ bg
$\tau_\mu \tau_\mu$	3	large $Z/\gamma^* \rightarrow \mu\mu$ bg
$\tau_e \tau_\mu$	6	low jet backgrounds
$\tau_e \tau_{had}$	23	golden
$\tau_\mu \tau_{had}$	23	golden
$\tau_{had} \tau_{had}$	41	challenging (jet bg's)

- Look for inclusive production of $A \rightarrow \tau\tau$:
 - Select di-tau candidate events: $\tau_e \tau_\mu, \tau_e \tau_{had}, \tau_\mu \tau_{had}$
 - No requirement on associated particle production
 - Construct a variable that gives Z/Higgs separation (m_{vis})
 - Extract signal/set limits by fitting the observed distribution with background and Higgs signal templates (various masses)

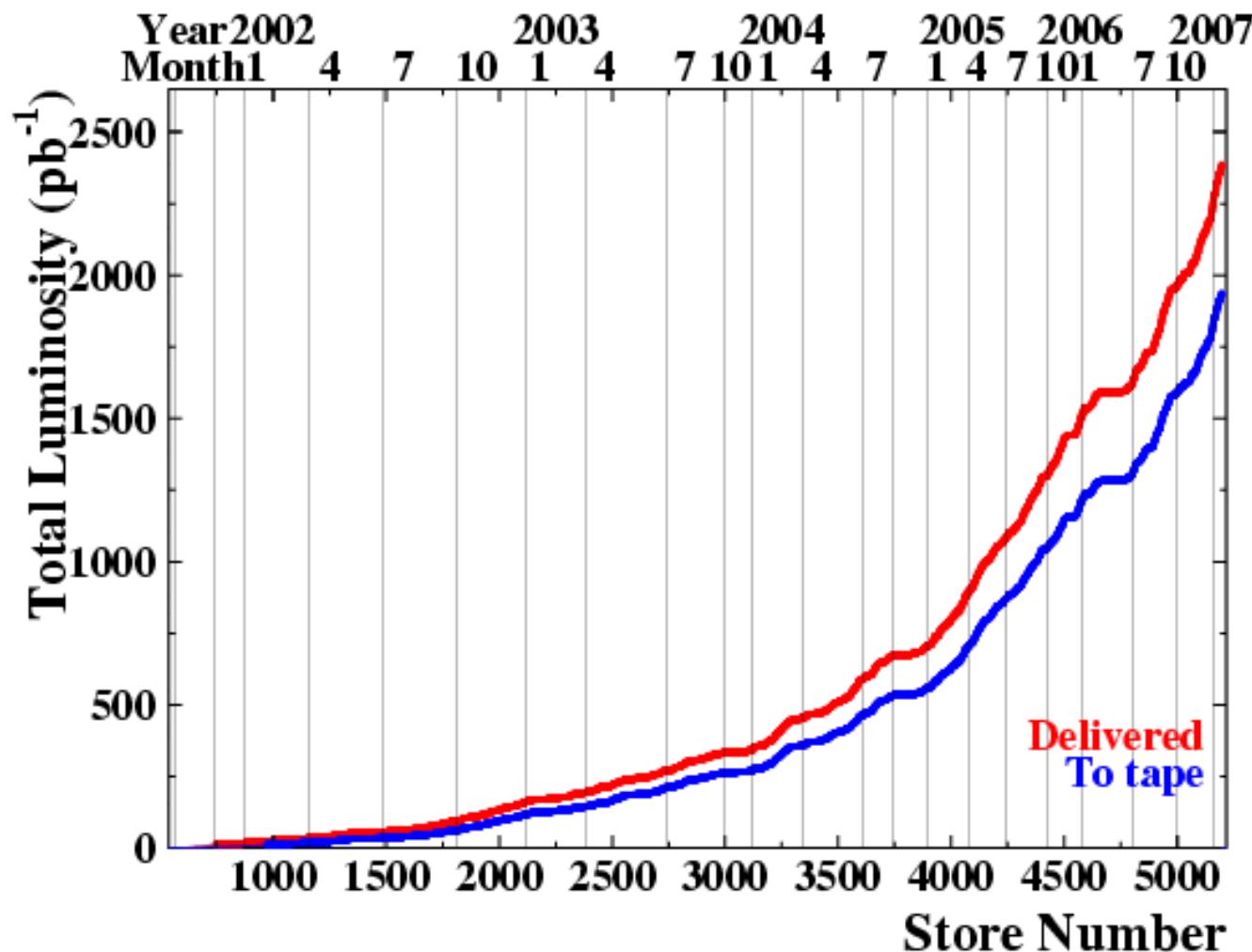
The Fermilab Tevatron



The CDF Detector at the Tevatron



- Tracking: Silicon detectors and wire chambers
- Calorimetry: EM, Hadronic
- **EM Shower Maximum detector** - used in π^0 reconstruction
- Muon Detectors
- TOF System



CDF has collected almost 2 fb^{-1} of data (Jan 2007)!
This analysis uses 1 fb^{-1} .

- Dedicated set of tau triggers in Run 2:

- Electron + τ_{had}**

- Central electron ($E_T > 8 \text{ GeV}$) + isolated track ($p_T > 5 \text{ GeV}$)

- Muon + τ_{had}**

- Central muon ($p_T > 8 \text{ GeV}$) + isolated track ($p_T > 5 \text{ GeV}$)
no tracks with $p_T > 1.5 \text{ GeV}$
in isolation annulus $10-30^\circ$

- Missing Transverse Energy + τ_h**

- $\cancel{E}_T > 20 \text{ GeV}$ + narrow isolated jet

- Di-tau: $\tau_{\text{had}} + \tau_{\text{had}}$**

- Two narrow, isolated jets
Cluster-matched track,
no tracks with $p_T > 1.5 \text{ GeV}$ in $\Delta\phi = 10-30^\circ$

- Low- p_T di-lepton triggers

- $e\bar{e}/\mu\bar{\mu}/\bar{e}\mu$; $p_T > 4 \text{ GeV}$**

Electron and muon selection

Electrons

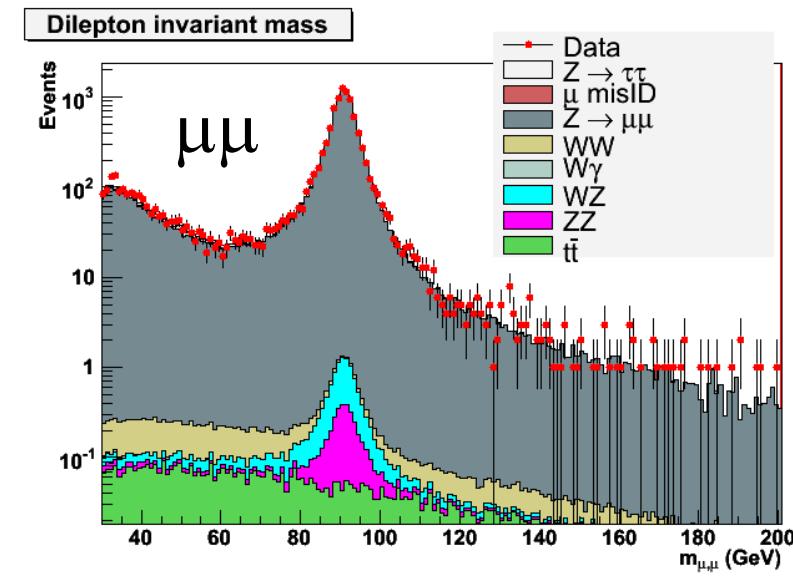
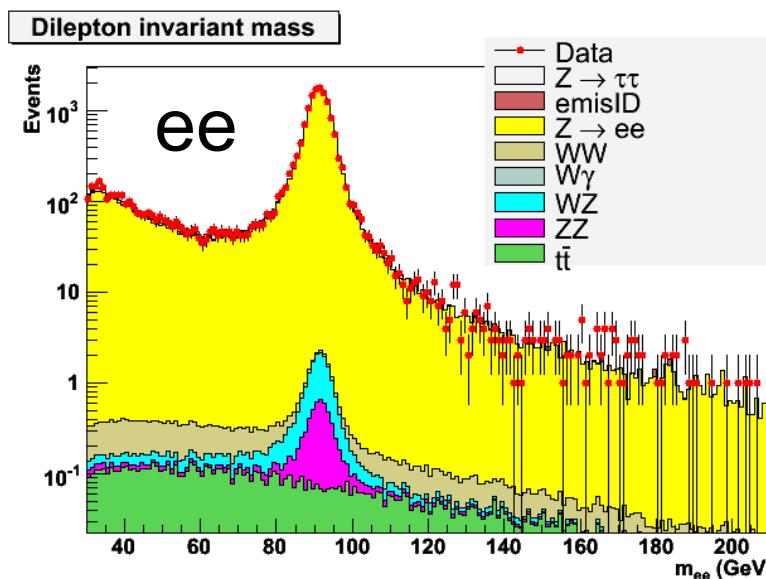
- In the central detector
- $E_{\text{HAD}}/E_{\text{EM}}$, E/p , and calorimeter cluster lateral shape consistent with the expectation for electrons
- Not part of a conversion pair
- $E_T > 10 \text{ GeV}$ (6 GeV) for $\tau_e \tau_{\text{had}}$ ($\tau_e \tau_\mu$)

Muons:

- Fiducial in the central muon detectors
- Hits in the muon detectors associated with the track
- Small energy deposition in the calorimeter
- $p_T > 10 \text{ GeV}$ (6 GeV) for $\tau_\mu \tau_{\text{had}}$ ($\tau_e \tau_\mu$)

Isolation (same for e/ μ) :

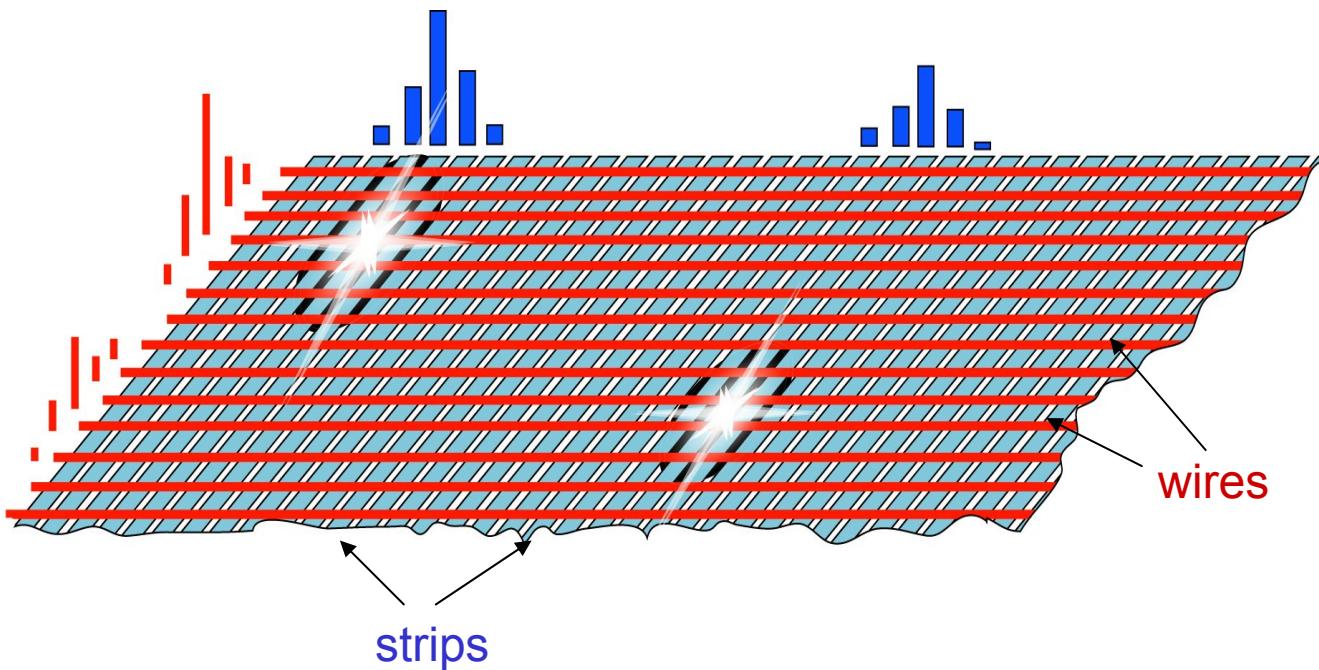
- $\sum p_T^{\text{trk iso}}(0.4) < 2 \text{ GeV}$ ($\tau_e \tau_\mu$ final state)
- $I^{\text{cal}}(0.4) < \max(2.0, 0.1 \times p_T(\text{or } E_T)) \text{ GeV}$ (additional requirement for $\tau_{e/\mu} \tau_{\text{had}}$)



Throughout this presentation “tau reconstruction” refers to reconstruction of the hadronic system in $\tau \rightarrow \text{hadrons } v_\tau$

- Detector signature
 - Narrow jet
 - Low tracks and neutrals multiplicity
 - 1 or 3 trks, most of the time $\leq 3\pi^0$'s
 - 1-prong : 3-prong hadronic decays $\sim 3:1$
 - Most of the time the tracks are π^\pm (small fraction of K^\pm)
 - Mass of had system $m_{\text{had}} < m_\tau$
- Tau ID requires detection of tracks and π^0 's

- Position from the ShowerMax detector (inside EM cal)
- Energy determined from EM energy in the tower, corrected for tracks hitting the tower



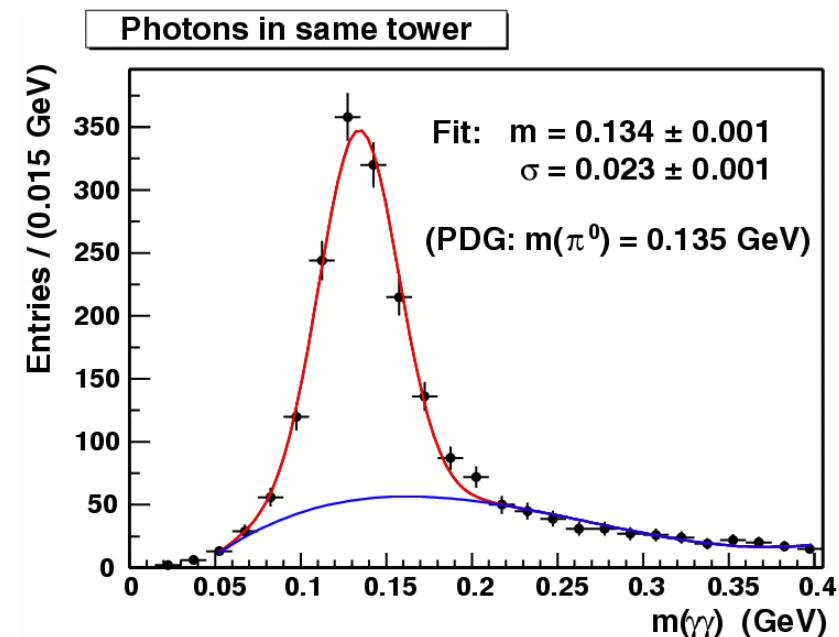
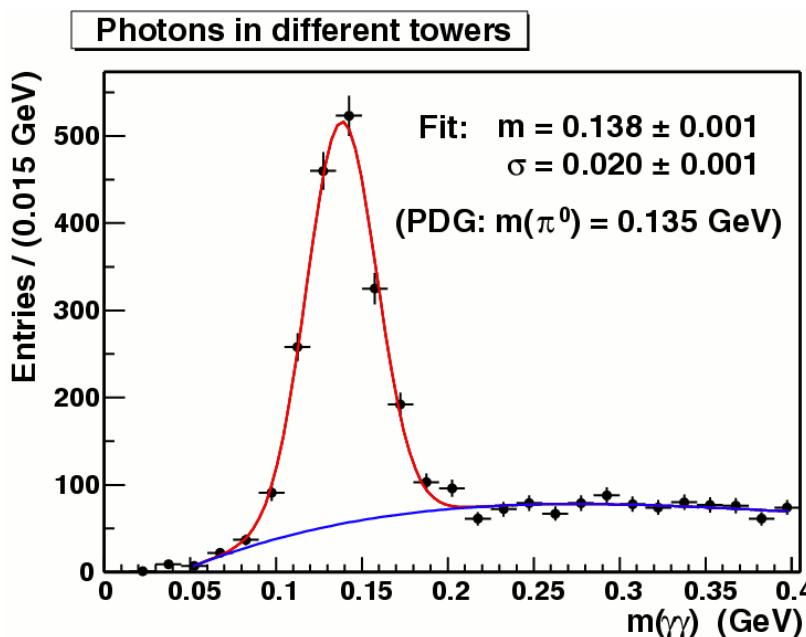
Orthogonal strips and wires,
independent strip/wire clusters

Each ShowerMax “module”
covers several towers

“Match” the strip to wire clusters
to obtain 2D positions of hits

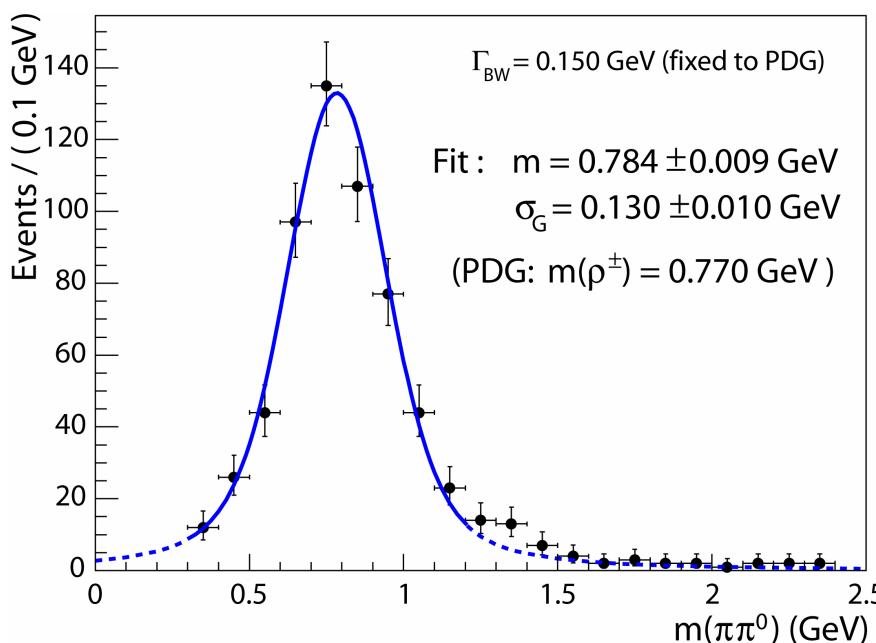
- Find CES matched clusters (not associated with tracks)
- Assign EM energy from the calorimeter tower:
If >1 match in the tower—split E_{EM} according to CES energy of each match

Well isolated, low- E π^0 's can be reconstructed from $\gamma\gamma$ (if we really try).



- Find CES matched clusters (not associated with tracks)
- Assign EM energy from the calorimeter tower:
If >1 match in the tower— split E_{EM} according to CES energy of each match

Most of the time a π^0 produces a single CES “match”.



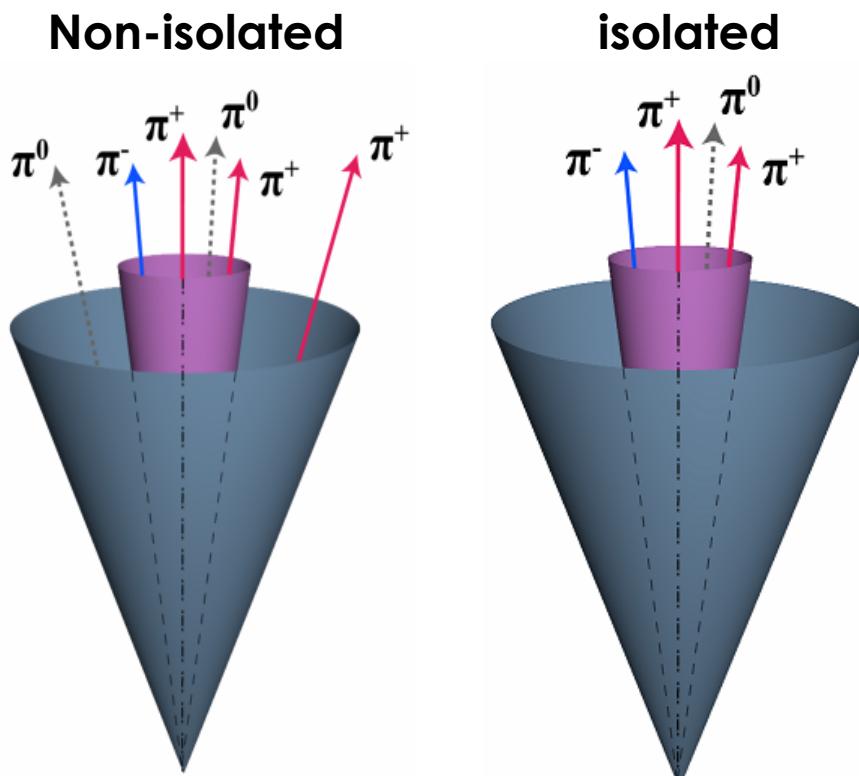
Select ρ 's from (mostly) $W \rightarrow \tau\nu$ decays
(illustration from previous studies)

- Exactly one trk and one π^0
- MET > 20 GeV (QCD clean-up)
- $p_T(\pi^0) > 6 \text{ GeV}$ (want merged γ 's)
- **no** other particles in 30 degrees
- Remove candidates consistent with $e + brem$

Two-cone algorithm for trks and π^0 's:

- Signal cone \Rightarrow reconstruct $\mathbf{p}_{\text{had}}(E, \vec{p})$
- Iso annulus \Rightarrow implement g/q jet veto
- Common axis: direction of a “high- p_T ” track (seed track), $p_T > 6 \text{ GeV}$
- Track points to a calorimeter cluster ($N^{\text{twr}} \leq 6$) with a seed tower with $E_T > 6 \text{ GeV}$

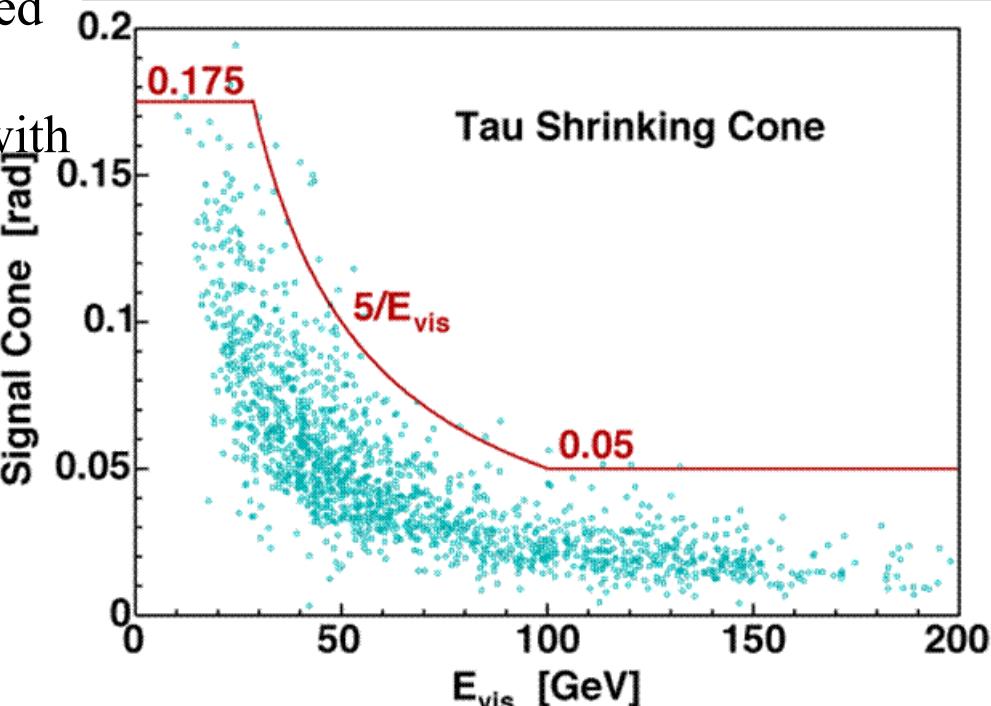
Tau candidates:



Signal cone size

$$\theta_{\text{sig}} = 5.0 \text{ (rad GeV)} / E_{\text{cl}} \text{ [rad]}$$

$$\theta_{\text{sig}}^{\max} = 0.174; \theta_{\text{sig}}^{\min}(\text{trk}) = 0.05; \theta_{\text{sig}}^{\min}(\pi^0) = 0.1$$



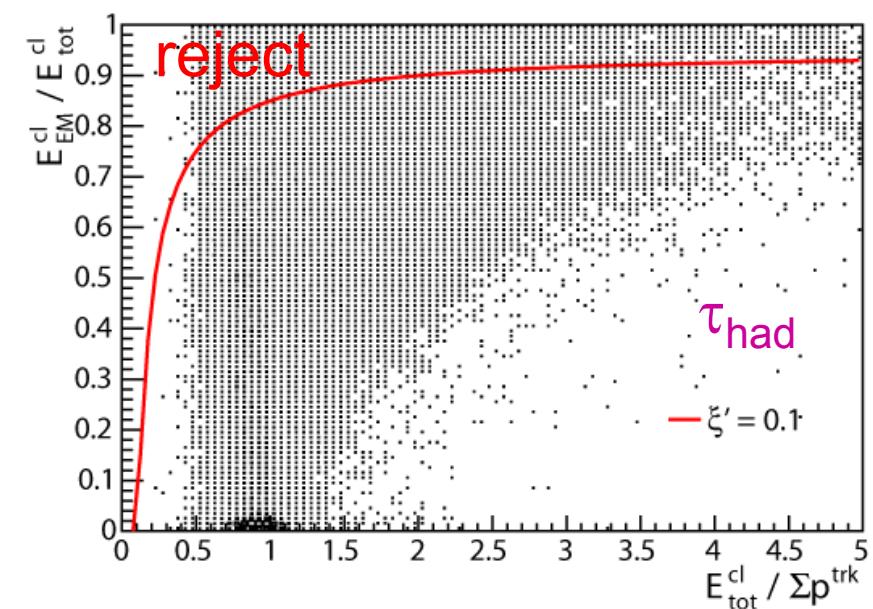
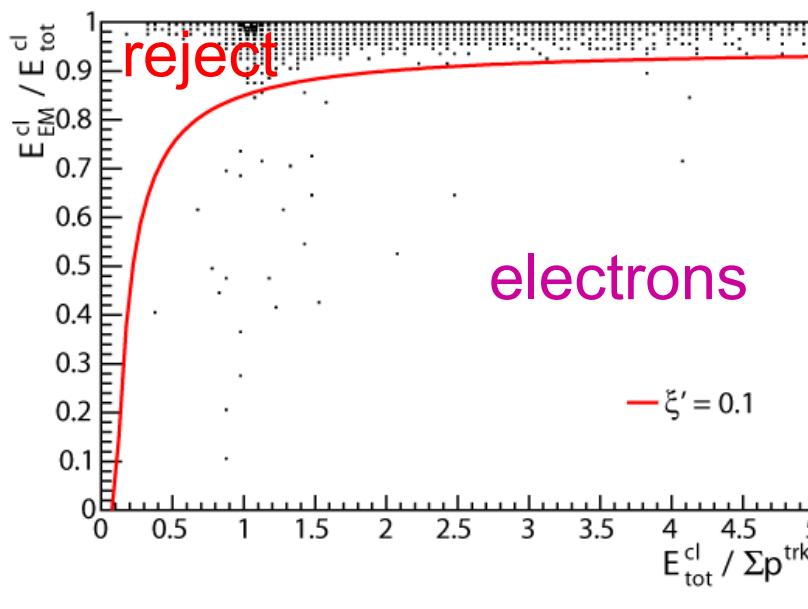
In tau reconstruction we assume that all tracks are π^\pm , and all reconstructed neutrals are π^0 's

Tau selection:

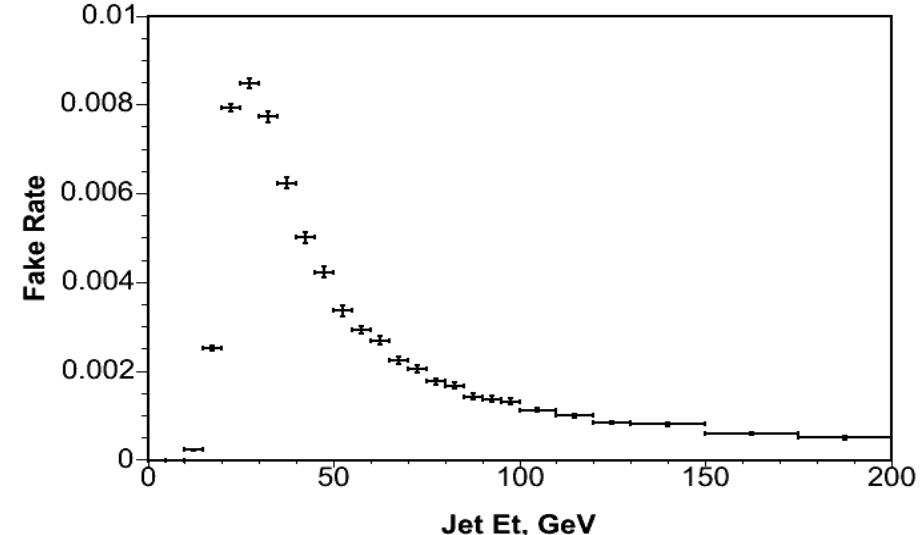
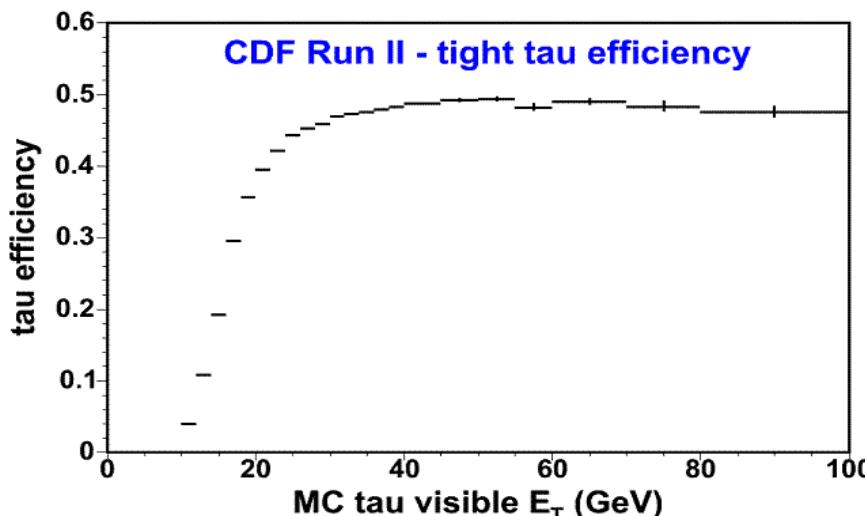
- number of tracks: $N_{\text{trk}} = 1, 3$
- $E_T^{\text{cl}} > 9 \text{ GeV}$,
- $p_T > 15 (20) \text{ GeV}$ for 1(3) prong
- $\xi' > 0.1$
- **track isolation:** $\sum_{\text{iso}} p_T^{\text{trk}} < 2 \text{ GeV}$
- π^0 isolation: $\sum_{\text{iso}} E_T^{\pi^0} < 1 \text{ GeV}$
- $m_{\text{had}} < 1.8 (2.2) \text{ GeV}$ for 1 (3) prong

$$\xi' = \frac{E_{\text{tot}}^{\text{cl}}}{\sum |\vec{p}_{\text{trk}}|} \left(0.95 - \frac{E_{EM}^{\text{cl}}}{E_{\text{tot}}^{\text{cl}}} \right)$$

Illustration of the effect of the ξ' cut on electrons and τ_{had}

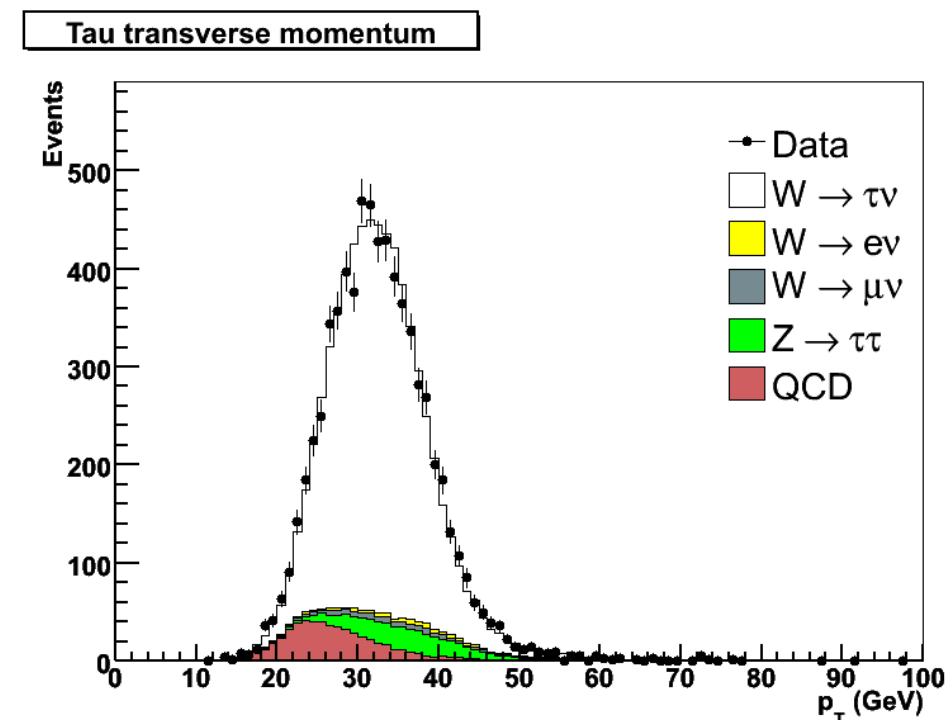


Tau efficiency and fake rate depend mostly on the isolation requirements



Data/MC energy scale

- Use taus from $W \rightarrow \tau\nu$ to compare the p_T spectra of τ_{had}
- Require one τ_{had} , $E_T > 25$ GeV, no extra jets in the event
- Compare the p_T spectra with MC shifted templates (KS tests)
- Data/MC scale consistent within $\sim 1\%$



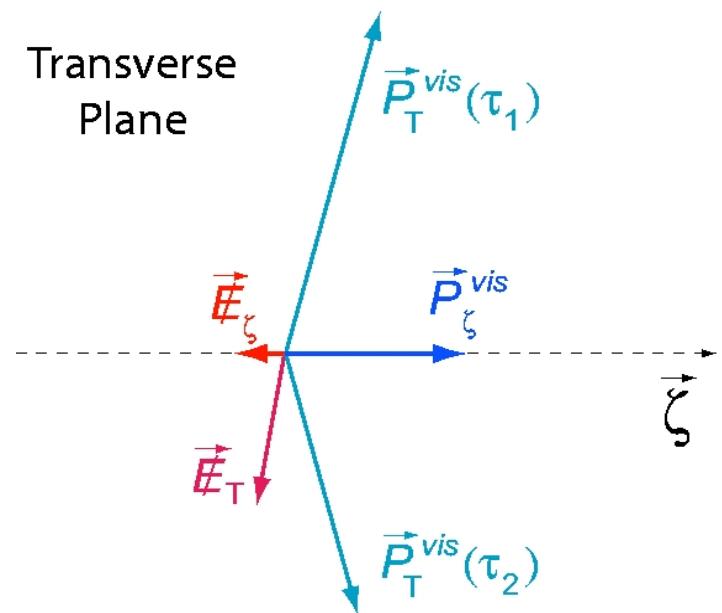
- $Z/\gamma^* \rightarrow \tau\tau$: estimated using MC
 - Largest background
 - Same final state as signal
- $Z \rightarrow \mu\mu, ee$: estimated using MC
- Misidentified jet as τ_{had}
 - multi-jet, W+jets; conversion+jet ($\tau_e \tau_{had}$ only)
 - Estimated from SS data events and W+jets MC
- Misidentified or non-isolated e/ μ in the $\tau_e \tau_\mu$ channel
 - Estimated from events in the e/ μ isolation sidebands
- “Other” : estimated using MC
 - ttbar, WW, WZ, ZZ, $W\gamma$, $Z\gamma$

Large number of background sources, most with small contributions

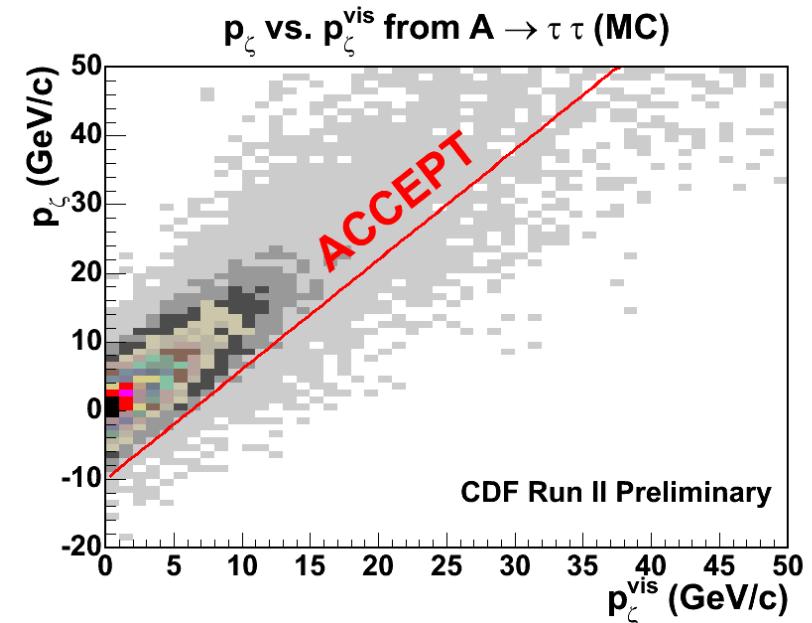
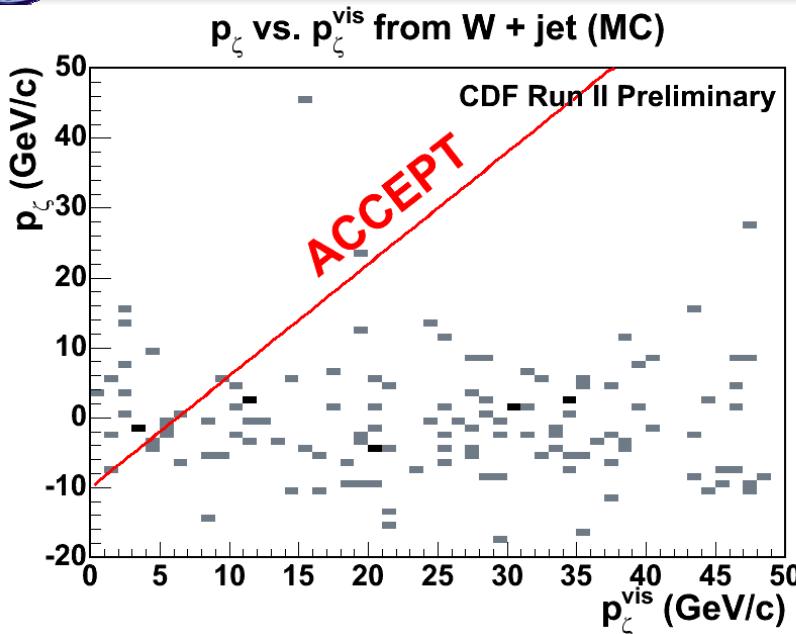
- $|z_0^{(1)} - z_0^{(2)}| < 5 \text{ cm}$
- $|0.5(z_0^{(1)} + z_0^{(2)})| < 60 \text{ cm}$
- $\Delta\phi(\tau^{(1)}, \tau^{(2)}) > 0.5 \text{ rad}$
- “Soft” QCD suppression
 - $\tau_{e/\mu} \tau_{\text{had}}$: $H_T = |\vec{p}_T| + |\vec{E}_T^e(p_T^\mu)| + |\vec{E}_T| >$ $\begin{cases} 55 \text{ GeV - 3-prong} \\ 45 \text{ (50) GeV - 1-prong } \tau_\mu \tau_{\text{had}} (\tau_e \tau_{\text{had}}) \end{cases}$
 - $\tau_e \tau_\mu$: $|\vec{E}_T^e| + |\vec{p}_T^\mu| > 30 \text{ GeV}$
- ζ cut : di-boson, tt, W+jet(s) suppression

$$p_\zeta > 1.6 \vec{p}_\zeta^{\text{vis}} - 10$$

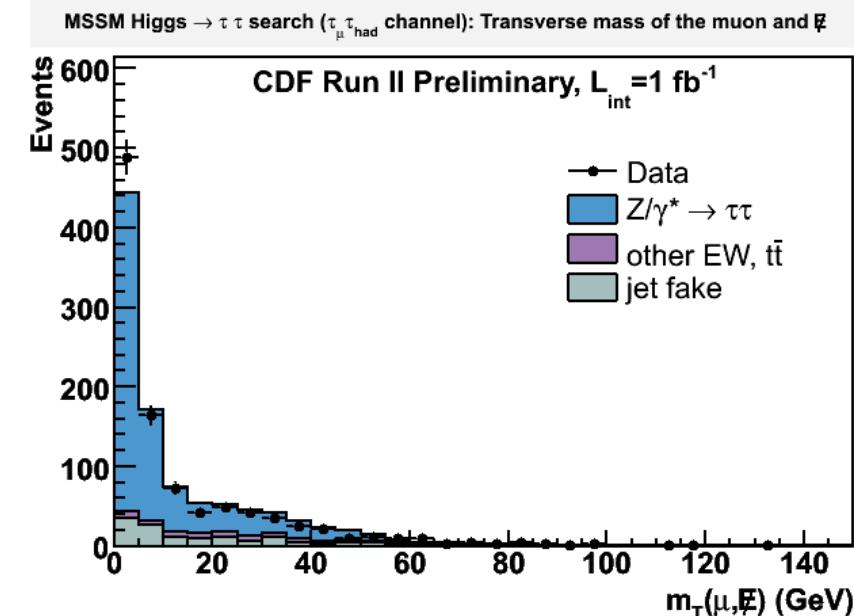
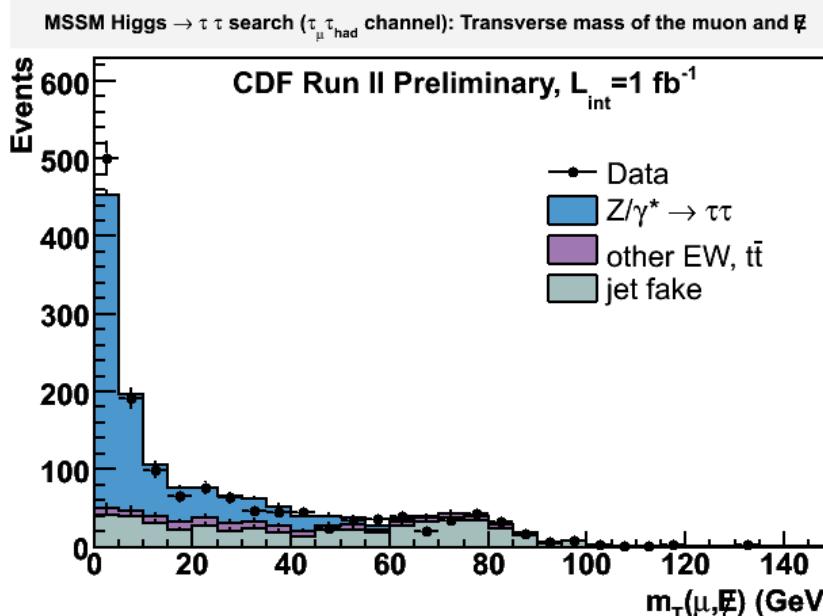
- ζ = bisection of $\Delta\phi(\tau_1, \tau_2)$
- $P_\zeta = (\vec{p}_T(\tau_1) + \vec{p}_T(\tau_2) + \vec{E}_T). \vec{\zeta}_n$
- $\vec{p}_\zeta^{\text{vis}} = (\vec{p}_T(\tau_1) + \vec{p}_T(\tau_2)). \vec{\zeta}_n$



Event Selection: the ζ -cut



Example: effect of the ζ -cut on $\tau_\mu\tau_{\text{had}}$ candidate events



Signal detection efficiency

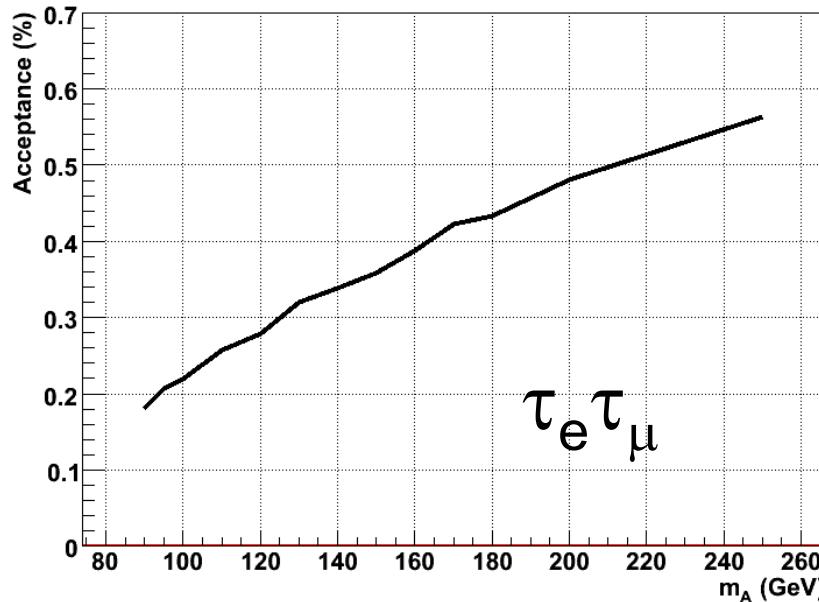
Signal MC generated with PYTHIA:

- $gg \rightarrow A$
- $bb \rightarrow A$

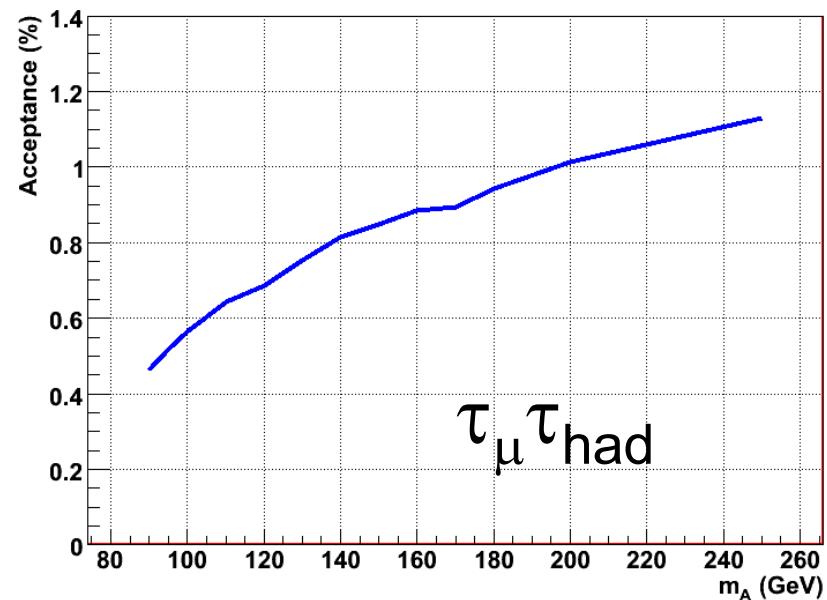
Higgs mass range: 90-250 GeV

Mass-dependent fractional contributions according to the expected cross sections

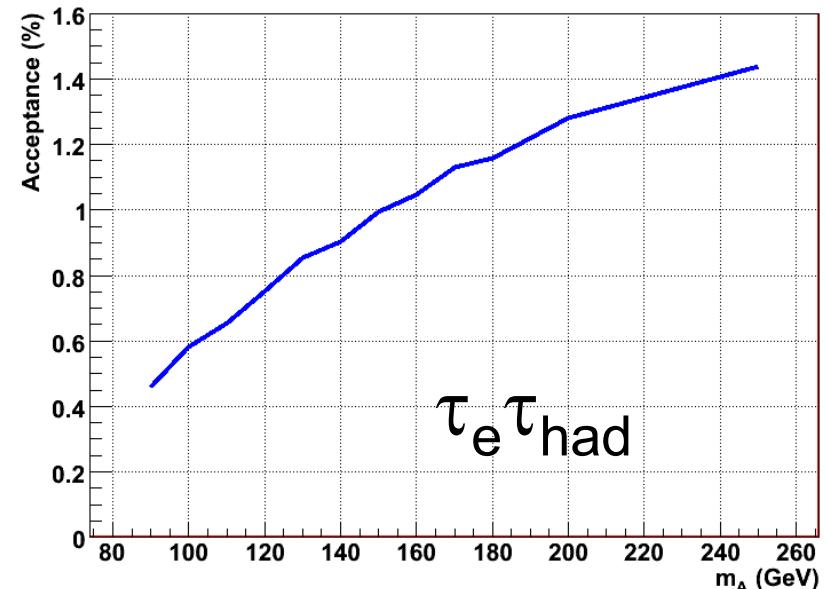
$A \rightarrow \tau_e \tau_\mu$ Total acceptance



$A \rightarrow \tau_\mu \tau_{had}$ Total acceptance

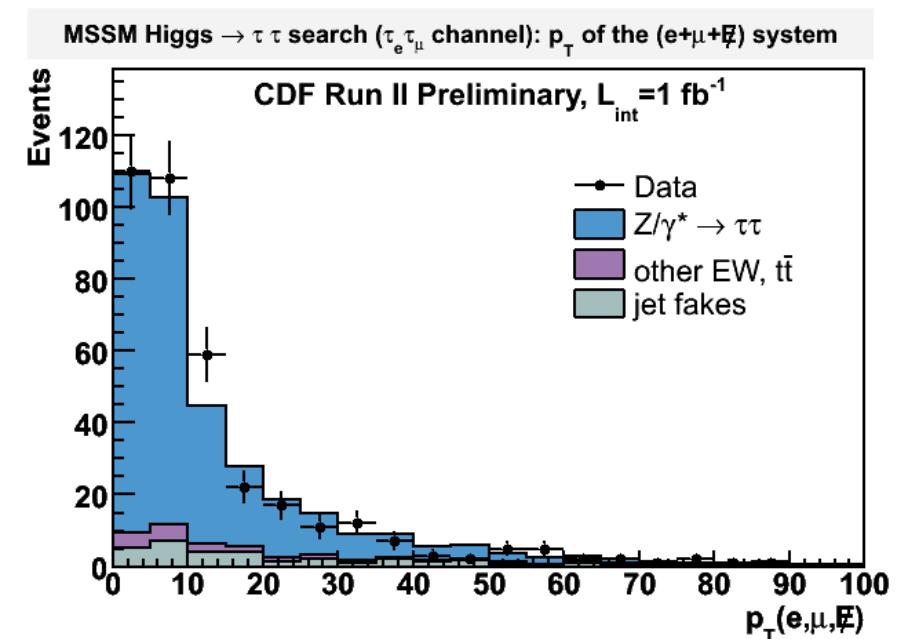
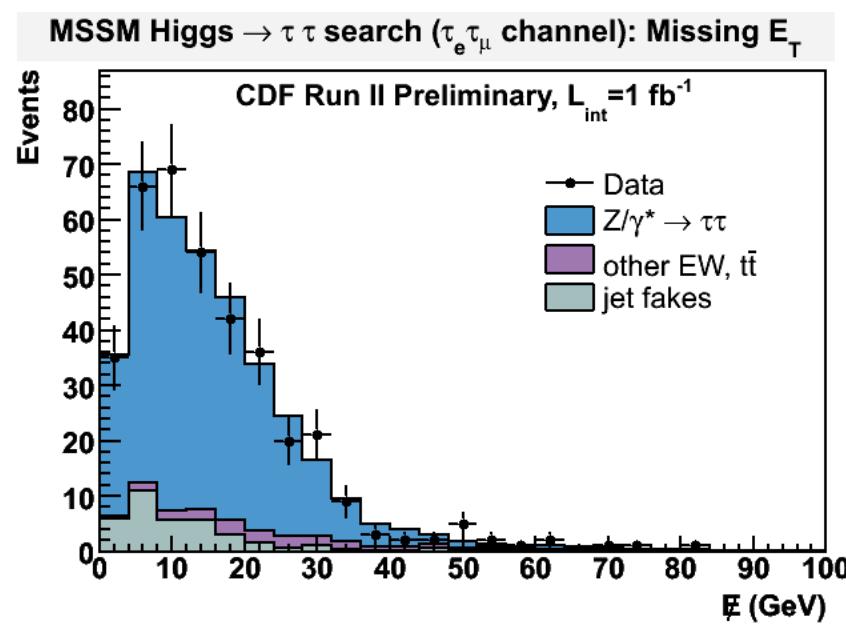
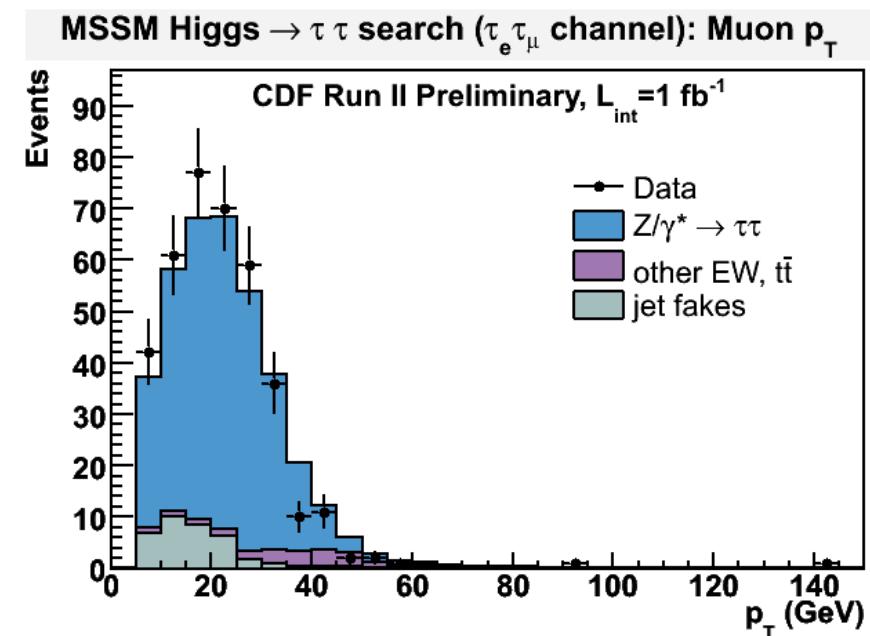
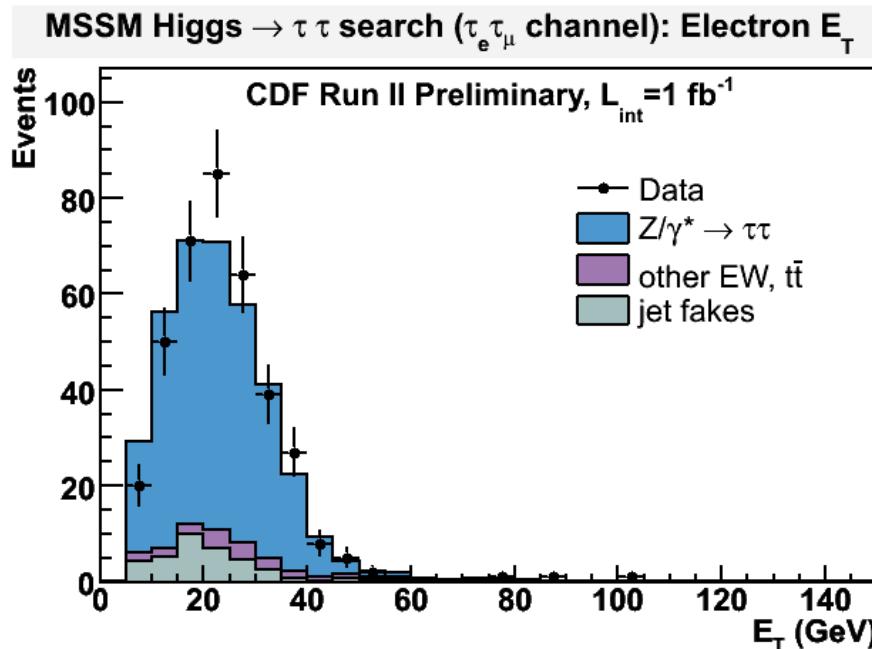


$A \rightarrow \tau_e \tau_{had}$ Total acceptance

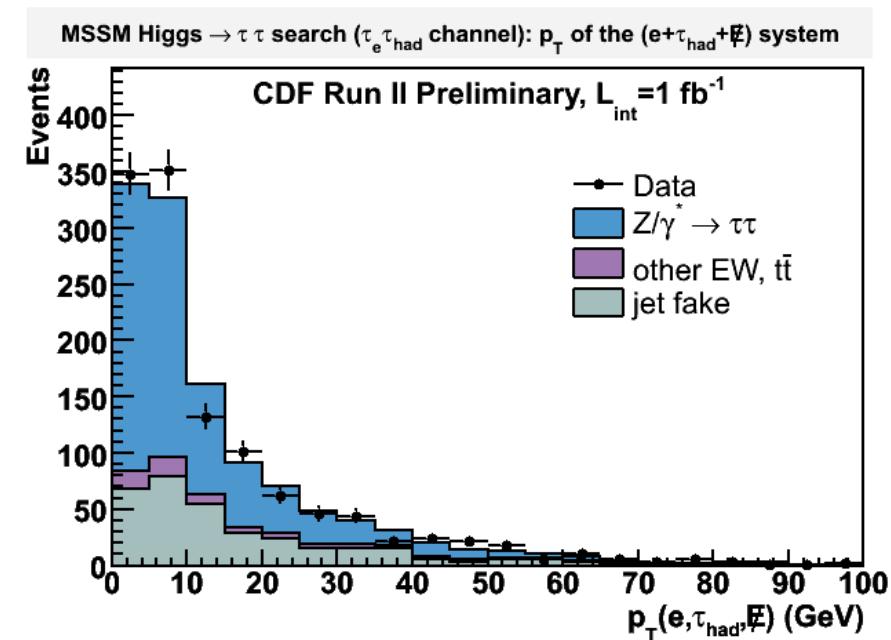
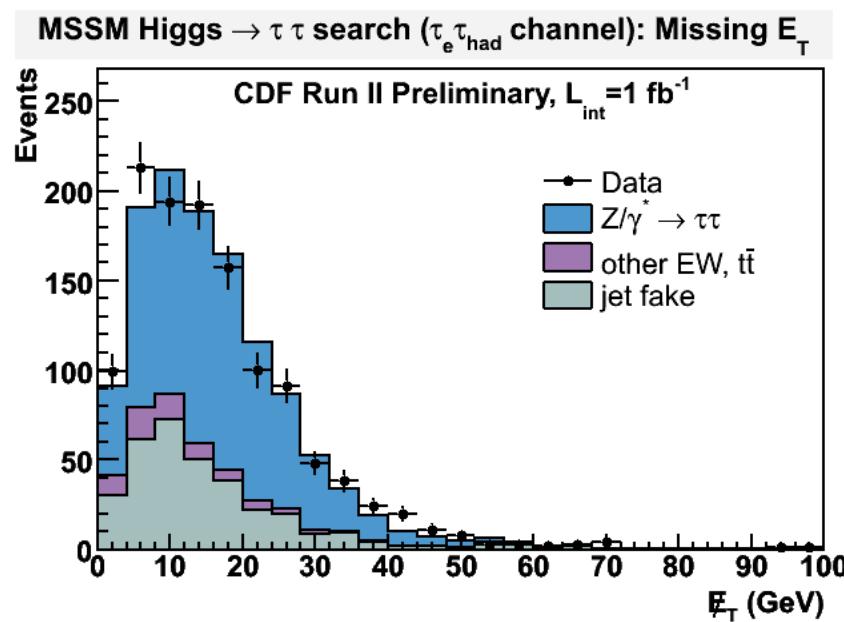
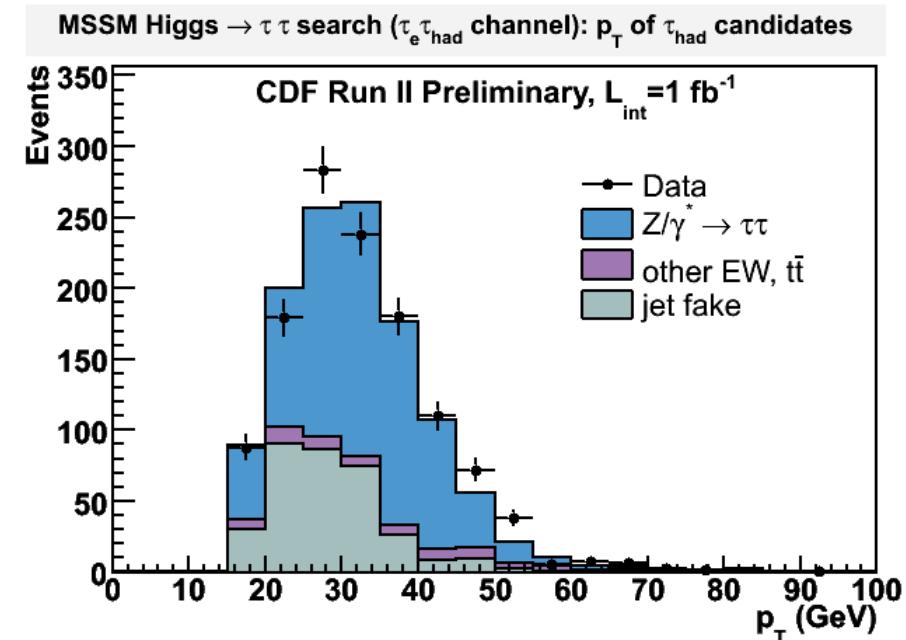
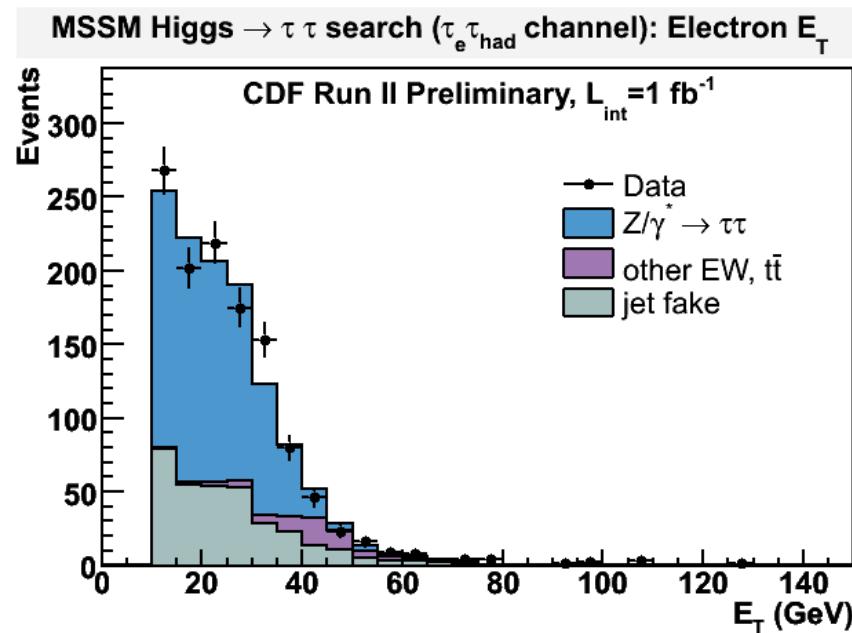


Tau BR's are included in the calculation of the acceptance. The total acceptance is the sum of the three numbers.

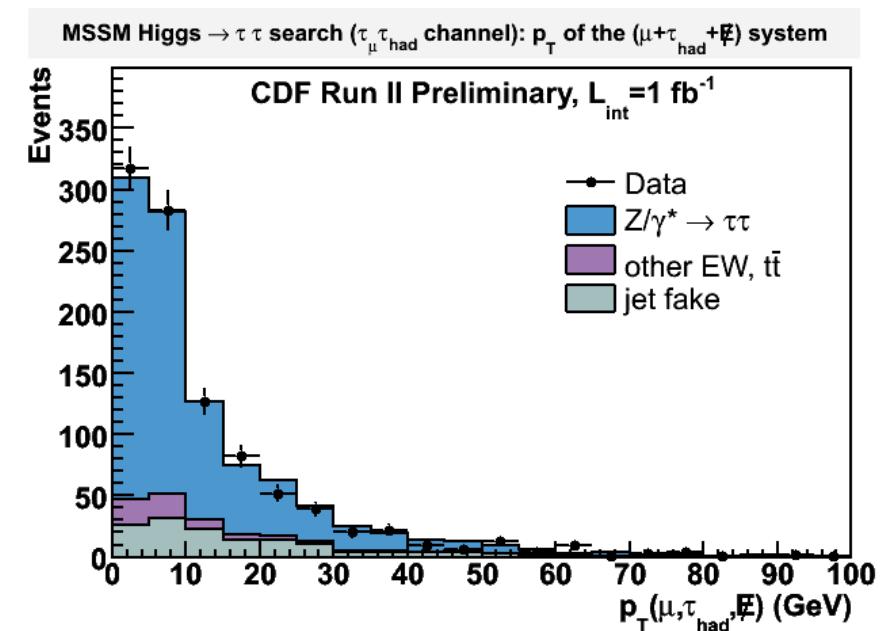
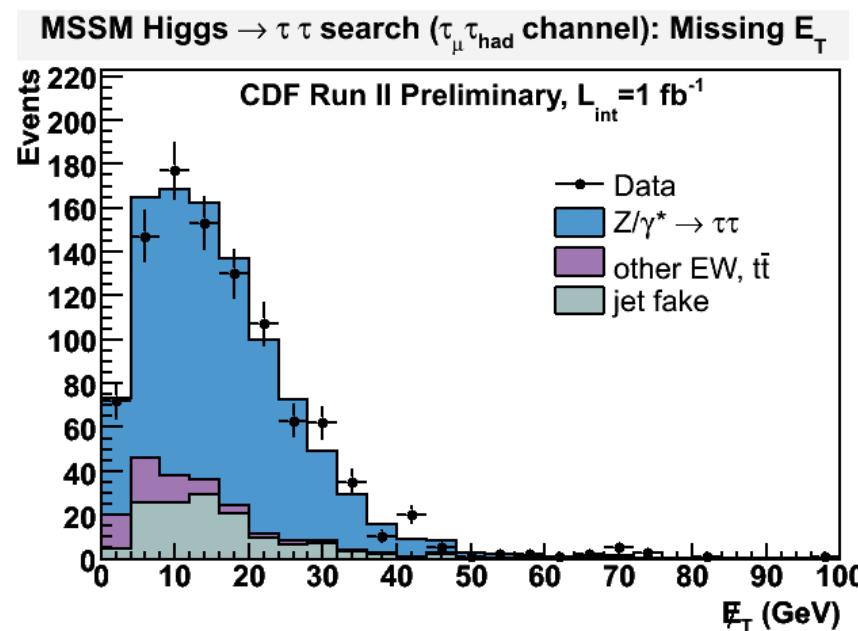
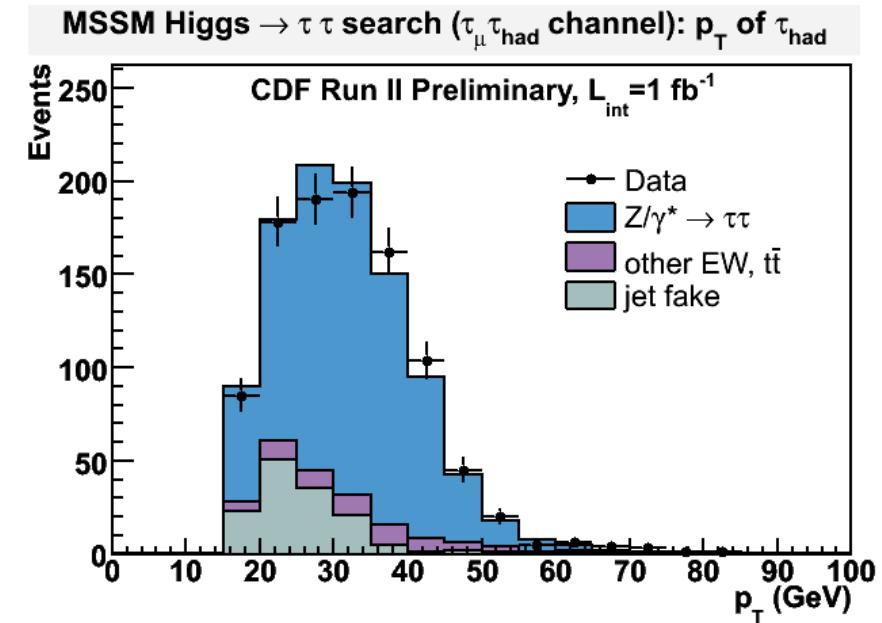
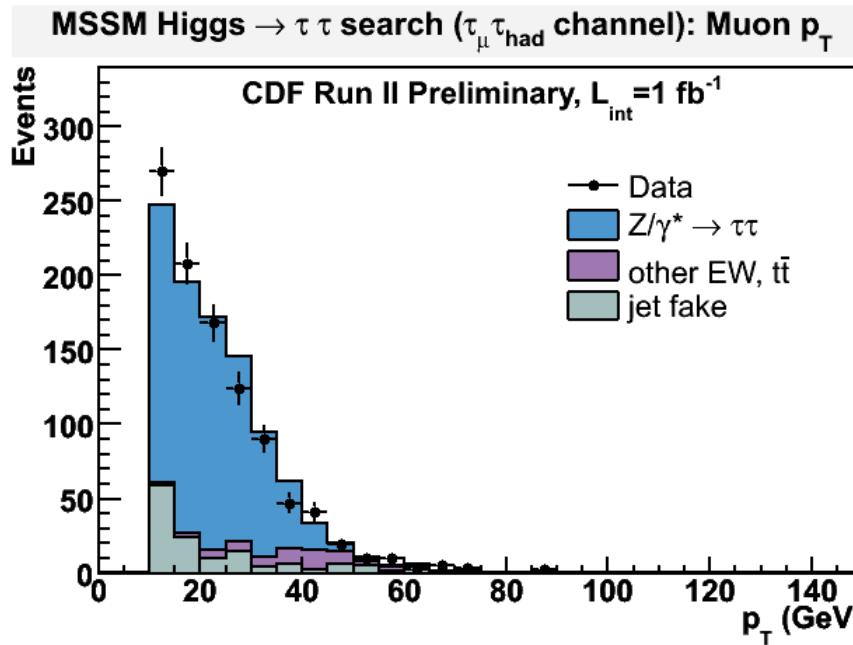
Selected events in the $\tau_e \tau_\mu$ channel



Selected events in the $\tau_e \tau_{\text{had}}$ channel



Selected events in the $\tau_\mu \tau_{\text{had}}$ channel



Observed events and expected backgrounds

Predicted bg's and observed events (stat errors only)

source	$\tau_e \tau_{had}$	$\tau_\mu \tau_{had}$	$\tau_e \tau_\mu$
$Z \rightarrow \tau\tau$	793.0 ± 4.7	796.6 ± 4.6	312.4 ± 2.9
$Z \rightarrow ee, \mu\mu$	68.3 ± 1.9	63.2 ± 1.8	11.9 ± 0.8
di-boson events	1.5 ± 0.02	1.2 ± 0.02	6.1 ± 0.1
$t\bar{t}$	1.3 ± 0.03	1.1 ± 0.03	4.7 ± 0.07
jet fakes	331.7 ± 18.2	139.4 ± 11.4	33.5 ± 3.2
Sum BG	1195.9 ± 18.9	1001.5 ± 12.5	368.6 ± 4.4
DATA	1215	1000	374

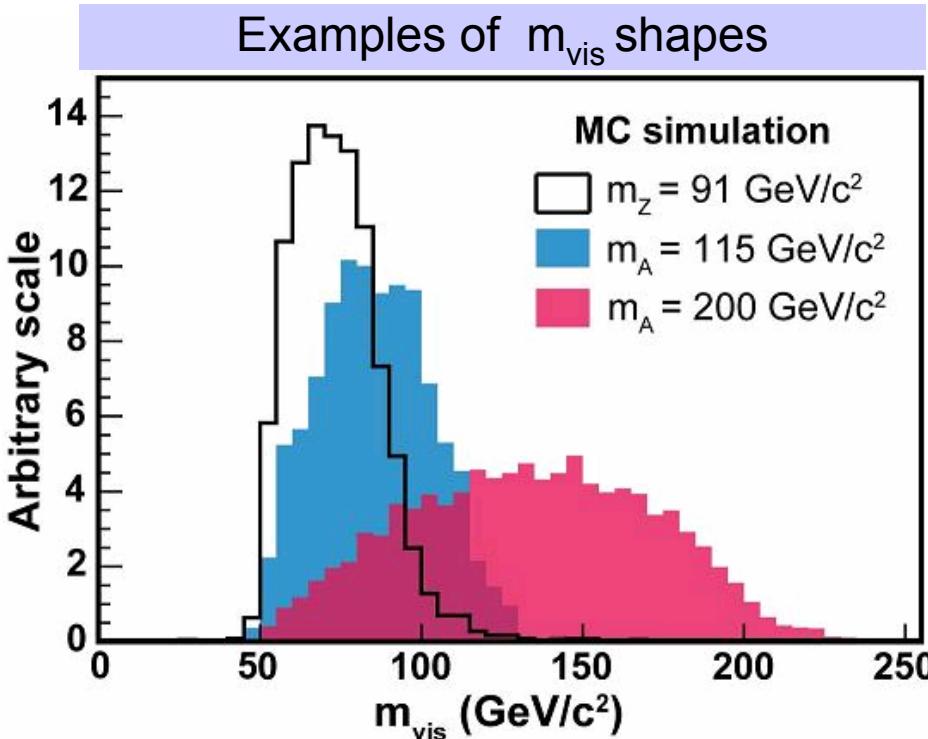
Two types of systematic uncertainties,
affecting:

- rates: data/MC detection efficiency
- distribution shapes: data/MC energy scales

Systematic uncertainties

Parameter	type	Error(%)	applies to:
e ID	rate	2.4	e in MC
muon reco+ID	rate	2.6	μ in MC
tau ID	rate	3.0	τ in MC
e trig	rate	0.3	e in MC
muon trig	rate	1.0	μ in MC
tau trig	rate	3.0	τ in MC
z-vertex cut	rate	0.5	all MC
fake/non-iso lepton bg in $\tau_e \tau_\mu$	rate	32.0	fake/non-iso bg
fake τ_{had} in $\tau_e \tau_{had}$	rate	6.0	fake tau bg
fake τ_{had} in $\tau_\mu \tau_{had}$	rate	9.0	fake tau bg
$\sigma \times \mathcal{B}(Z \rightarrow ll)$	rate	2.2	Z MC
$\sigma(t\bar{t})$	rate	13.4	$t\bar{t}$ MC
di-boson cross sections	rate	10	di-boson MC
PDFs (<i>Higgs</i>)	rate	5.7	signal
Luminosity	rate	6.0	all MC
JES	shape	$\pm 1\sigma$ (per jet)	all MC
EM scale	shape	± 1.0	e in MC
Tau p_T scale	shape	± 1.0	τ_{had} in MC

- Not enough information for exact $m_{\tau\tau}$ reconstruction:
 - Approximation: project E_T onto vis τ decay products (see next slide)
 - but... significant loss of statistics
 - compromise: use partially reconstructed mass $m_{vis} = m(\tau_1^{vis}, \tau_2^{vis}, E)$



Partial mass:

- Weaker discriminator than full mass
 - Non-linear relation with mass
- ✓ Still, there is substantial non-overlapping region
- ✓ Can be calculated for all events

A look at fully reconstructed $\tau\tau$ masses

At this point we do not use the fully reconstructed mass to look for signal.

However, it is instructive to look at the distributions:

- for demonstrate what we expect in the future
- as a consistency check of event reconstruction

- Using projection of MET onto the direction of visible tau decay products (collinear approximation) one can reconstruct the full $\tau\tau$ mass *
- The mass reconstruction method of choice for large statistics
 - Taus must not be back-to-back in the transverse plane
 - MET must be pointing between the visible decay products
- The energies carried by the neutrinos are determined from the system

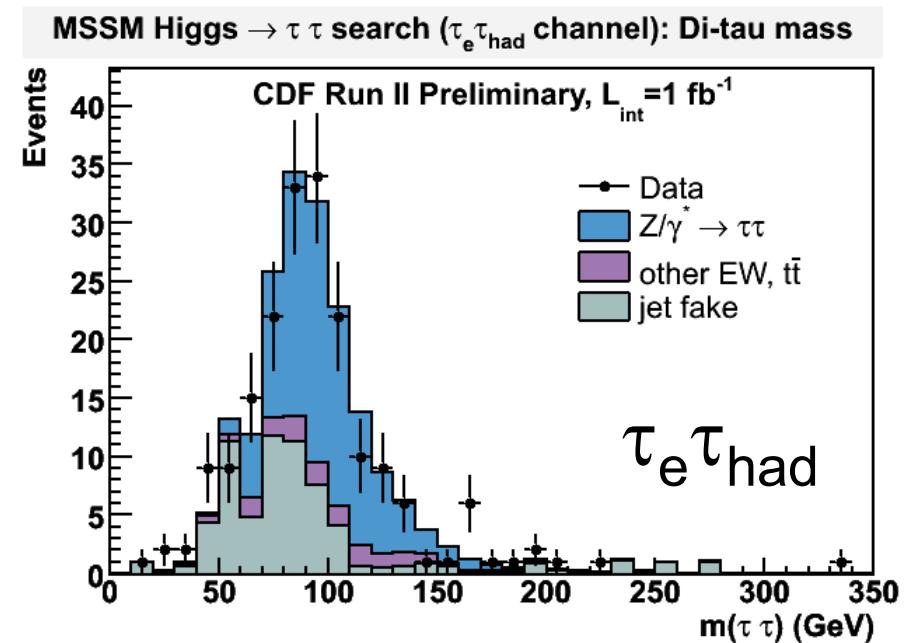
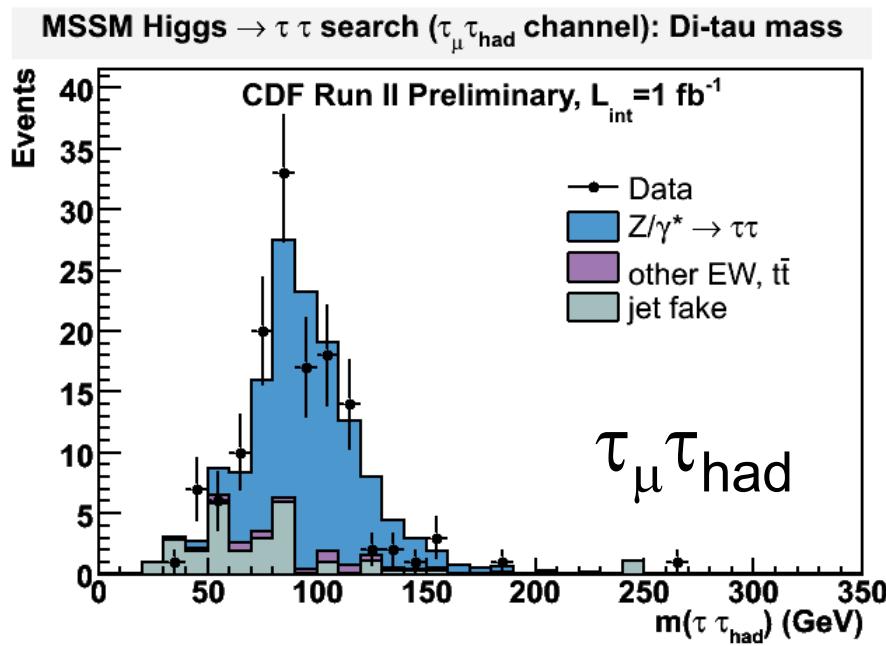
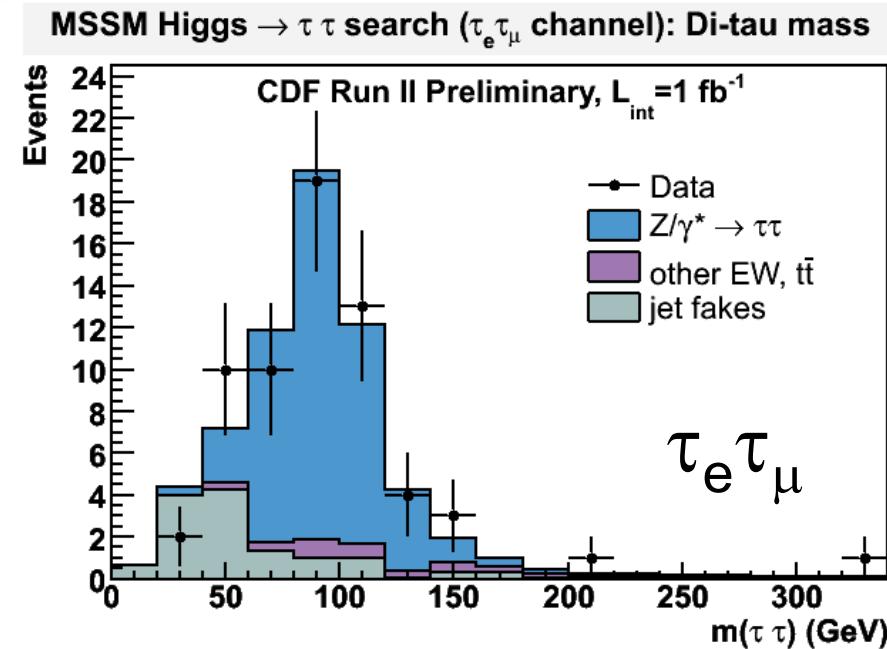
$$E_\nu^{(1)} \sin \theta^{(1)} \cos \phi^{(1)} + E_\nu^{(2)} \sin \theta^{(2)} \cos \phi^{(2)} = E_x$$

$$E_\nu^{(1)} \sin \theta^{(1)} \sin \phi^{(1)} + E_\nu^{(2)} \sin \theta^{(2)} \sin \phi^{(2)} = E_y$$

A look at fully reconstructed $\tau\tau$ masses

Non back-to-back taus: $|\sin(\Delta\phi)| > 0.3$

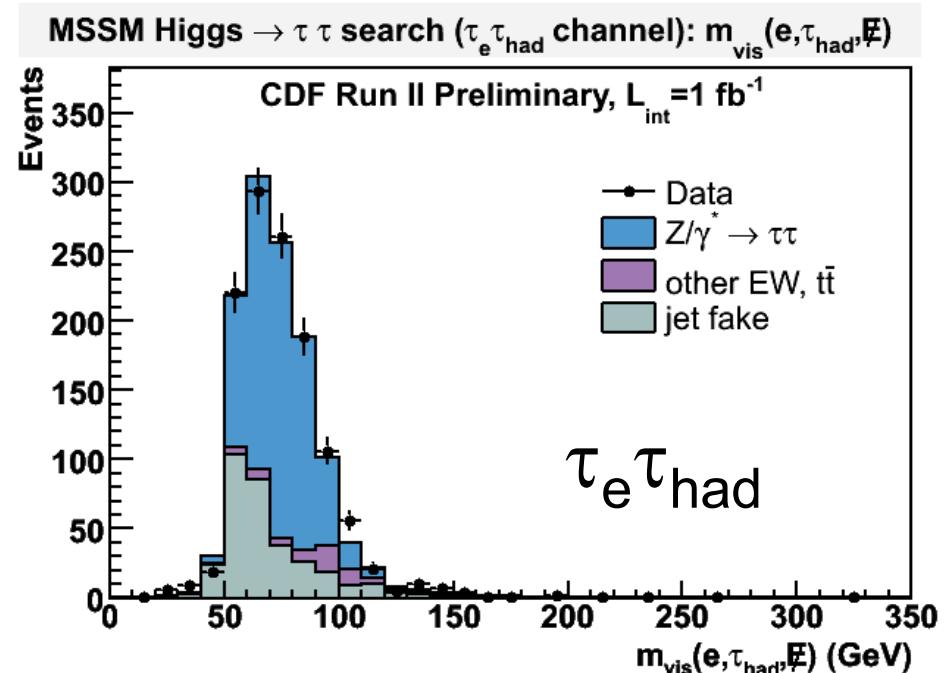
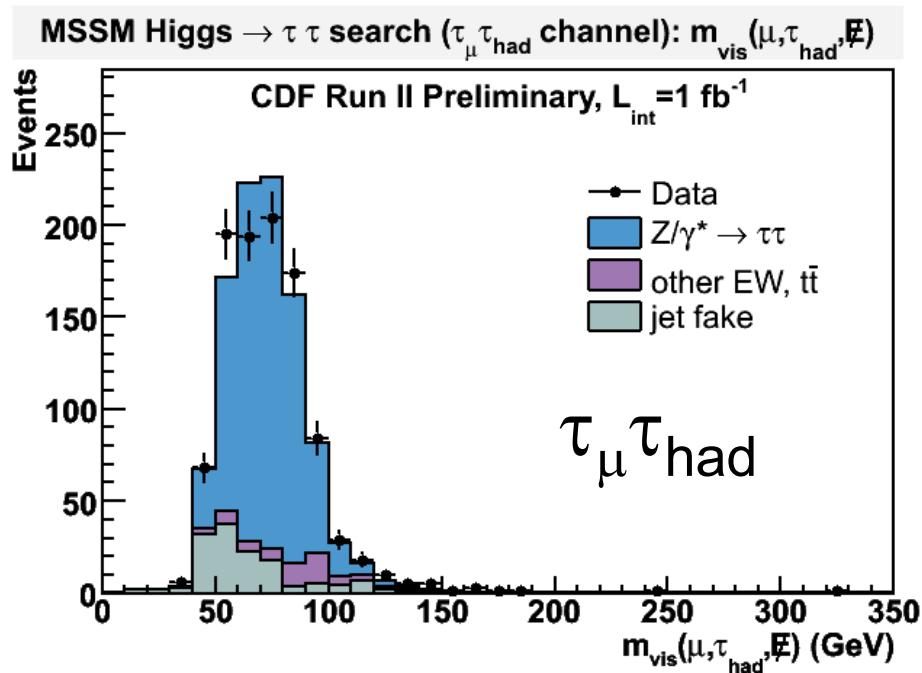
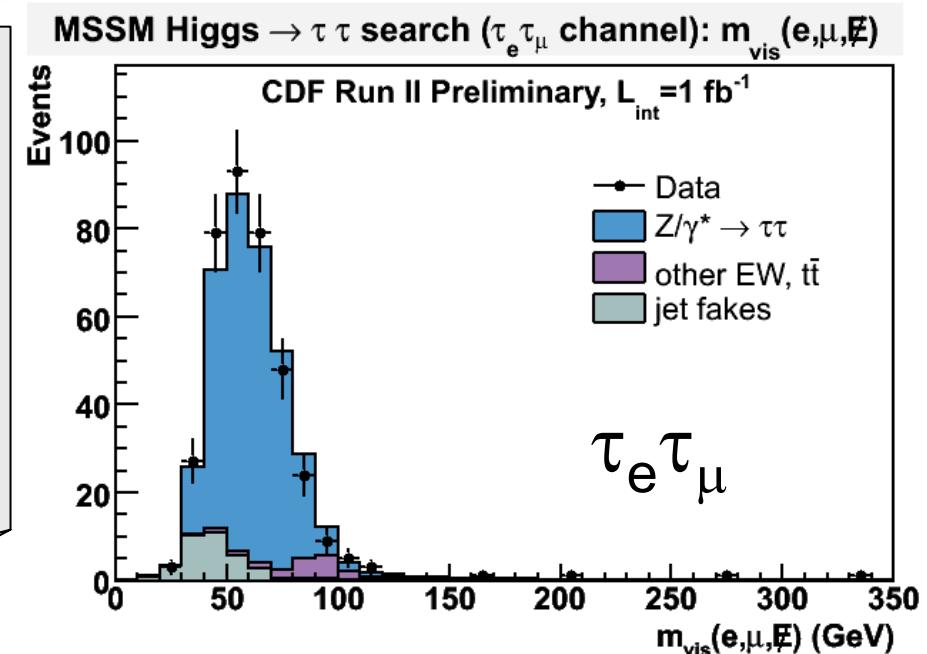
- the distributions peak at m_Z
- good agreement between data/MC
- significant loss of statistics

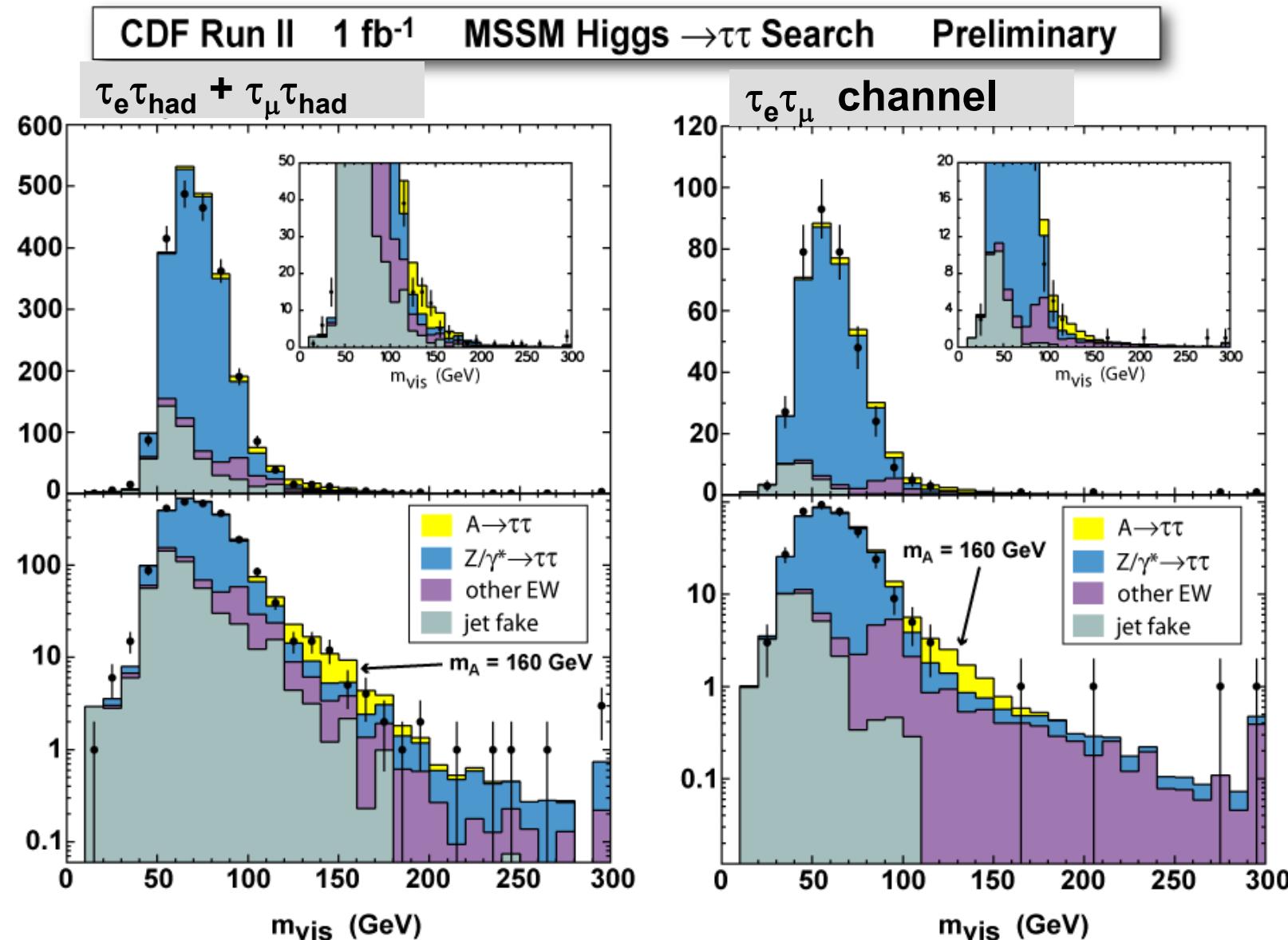


Limit extraction: from m_{vis} distributions

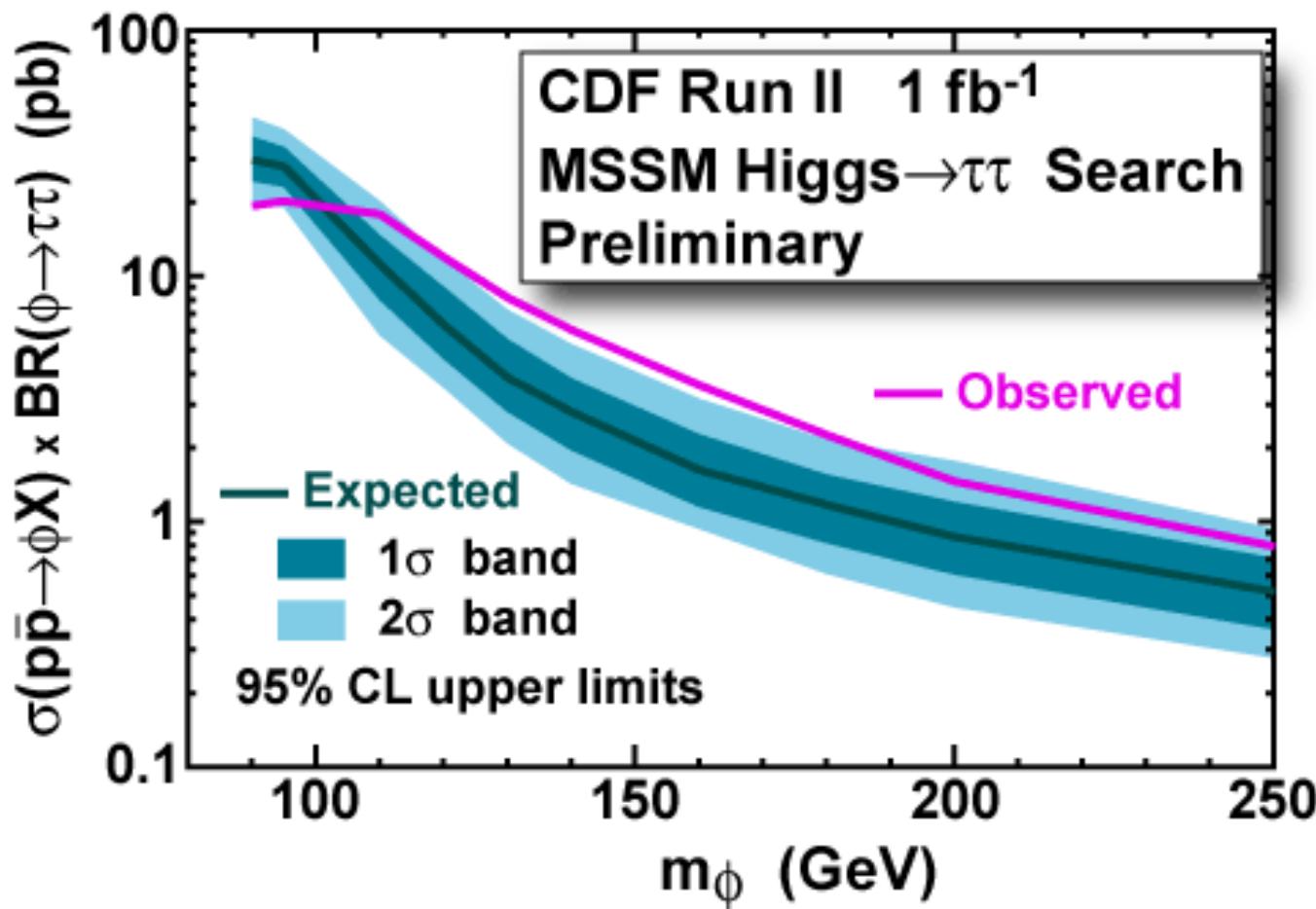
Likelihood fit to set limits:

- Treat the three final states as separate channels
- Backgrounds are allowed to float subject to Gaussian constraints
- correlations between backgrounds/channels are taken into account
- shape uncertainties included via template morphing





The backgrounds are normalized according to the fit result. The amount of Higgs signal corresponds to the exclusion at 95% CL.



- The excess of events results in weaker limits relative to the expectation, especially for $m_A=150-160$ GeV
 - The broad spectrum of m_{vis} leads to propagation of the effect to a large mass range
- Taking into account that the excess could have shown anywhere in the mass range, the effect is less than 2 sigma.

Interpretation of the results in MSSM

- Large MSSM parameter space, consider four benchmark scenarios ⁽¹⁾:

	M_{SUSY}	μ	M_2	X_t^{OS}	m_{gluino}
m_h^{max}	1 TeV	$\pm 200 \text{ GeV}$	200 GeV	$2 M_{\text{SUSY}}$	$0.8 M_{\text{SUSY}}$
no-mixing	2 TeV	$\pm 200 \text{ GeV}$	200 GeV	0	$0.8 M_{\text{SUSY}}$

- Follow the prescription of *Carena et al.* to get $\tan\beta$ from the $\sigma \times \text{BR}$ results, using:
 - $\sigma_{\text{SM}}(\text{gg} \rightarrow \phi + X)$ from HIGLU ⁽²⁾
 - $\sigma_{\text{SM}}(\text{bb} \rightarrow \phi + X)$ ⁽³⁾
 - $\text{BR}(\phi \rightarrow \tau\tau)$ and MSSM enhancement factors from the FeynHiggs program ⁽⁴⁾

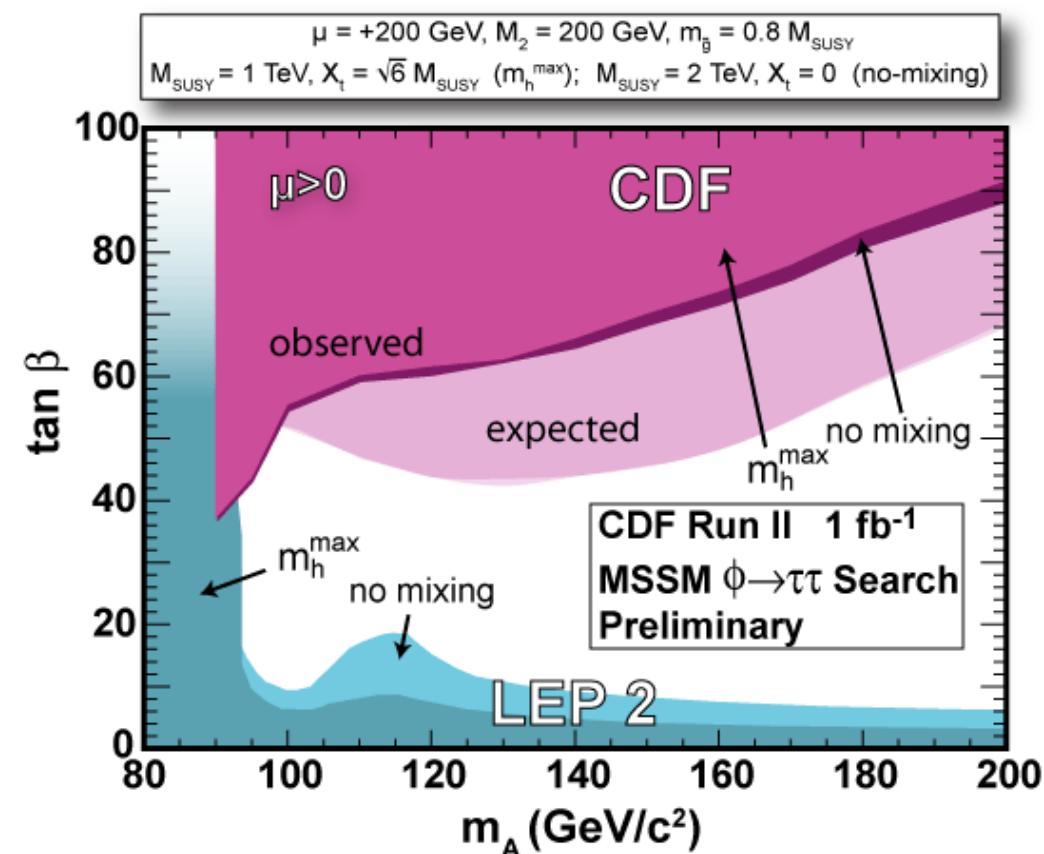
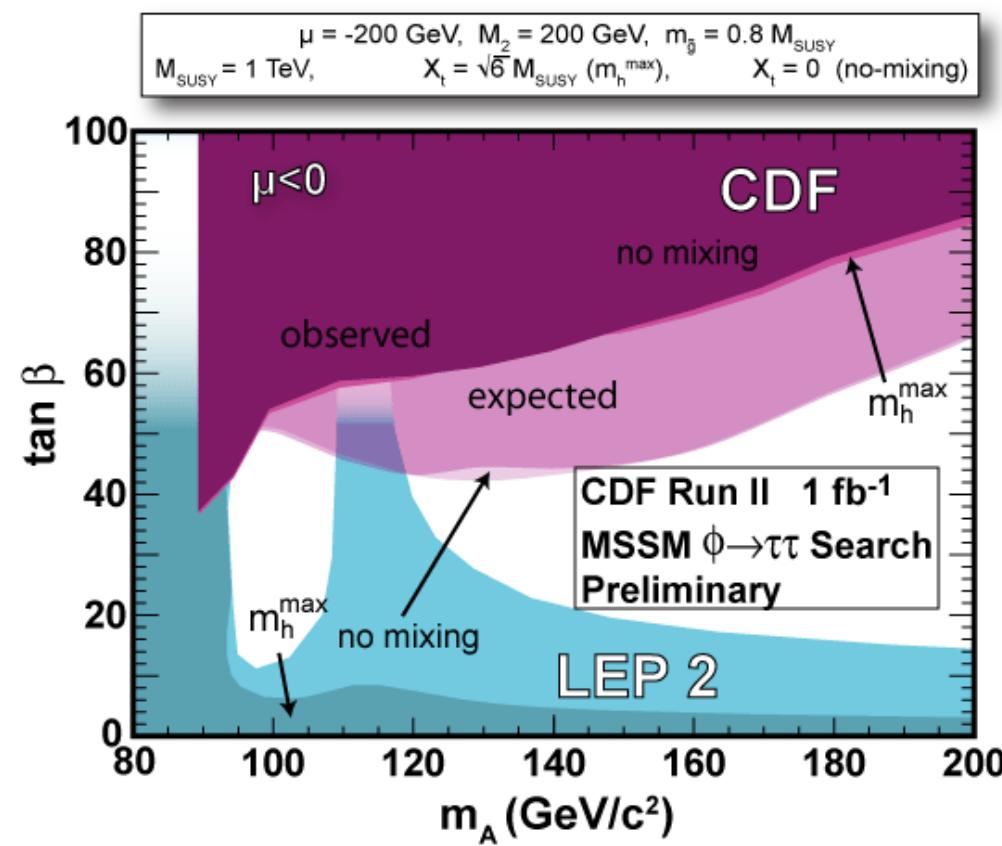
⁽¹⁾ M. Carena, S. Heinemeyer, G. Weiglein, and C.E.M. Wagner., Eur.Phys.J. C45 (2006) 797-814

⁽²⁾ M. Spira, Nucl. Instrum. Meth. A389, 357 (1997).

⁽³⁾ R.V. Harlander, W.B. Kilgore, Phys. Rev. D68, 013001 (2003).

⁽⁴⁾ S. Heinemeyer, W. Hollik, and G. Weiglein, Eur. Phys. J. C 9, 343 (1999); Comput. Phys. Commun. 124, 76 (2000).

$\tan\beta$ vs m_A exclusions

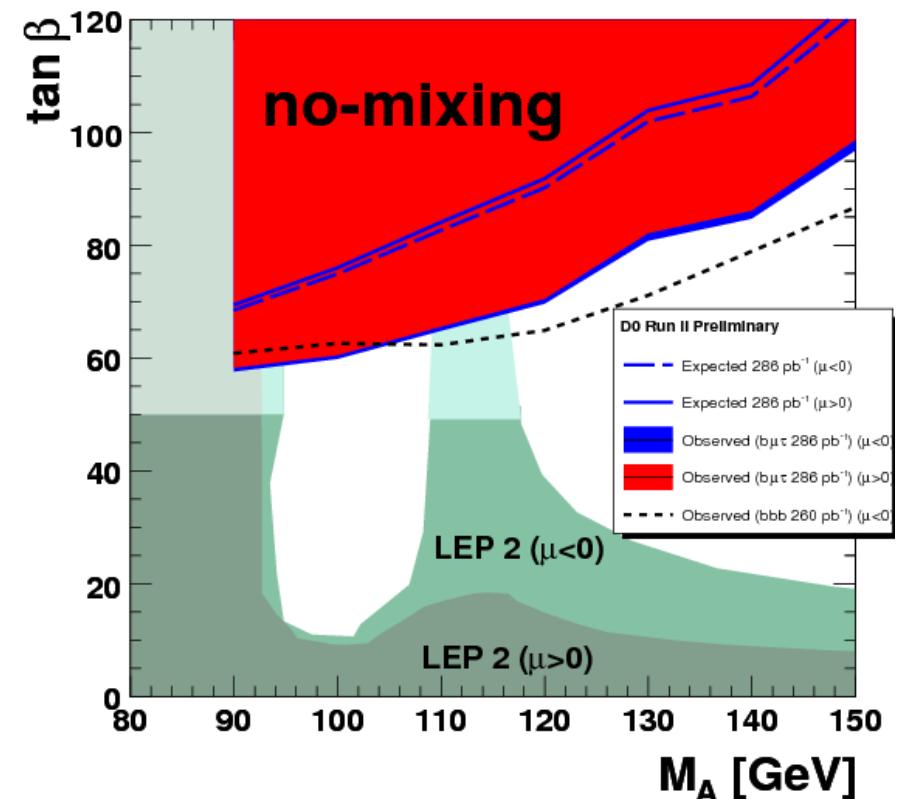
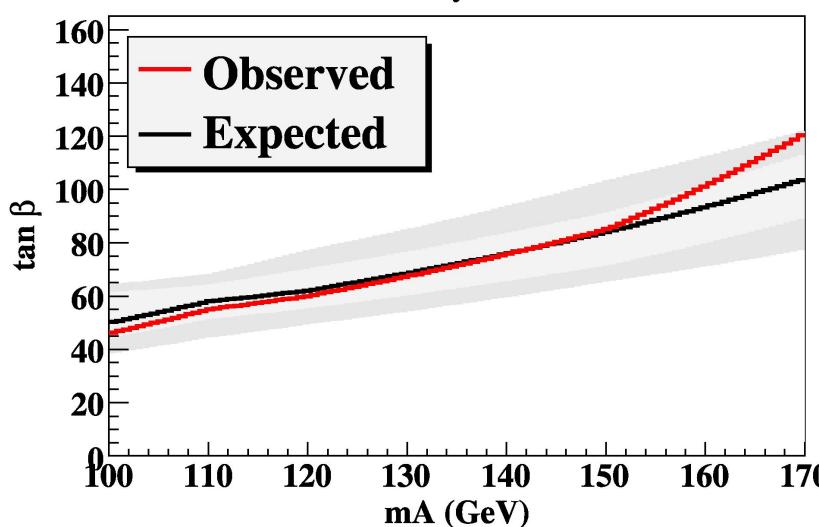
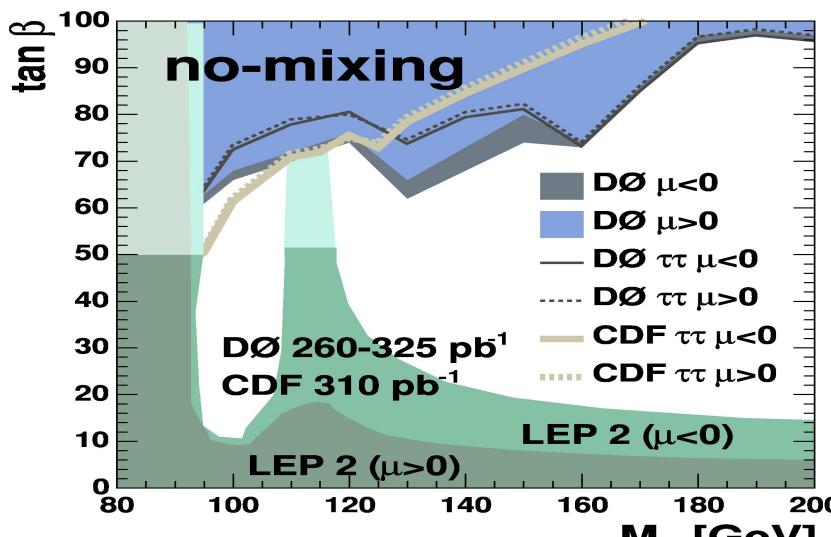


- As expected for the MSSM $H \rightarrow \tau\tau$ channel the exclusions in $\tan\beta$ vs m_A plane show small variation for the four scenarios

- We have observed a small, but interesting excess of events in the data
- CDF will soon have 2 fb^{-1} of data - a welcome and timely addition
- Refine the search utilizing more information, treat as separate channels events with
 - fully reconstructed $\tau\tau$ mass
 - associated high p_T b-jet

$\tan\beta$ vs m_A : other results from the Tevatron

- D0 has also performed searches in the $A \rightarrow \tau\tau$ (340 pb^{-1}), and $A+b(b) \rightarrow bbb(b)$ ($260, 900 \text{ pb}^{-1}$) modes
- New results from D0 and CDF will follow soon



It is no longer possible to overlay all the Tevatron results on a single plot



- The Tevatron experiments have a lot to say about MSSM Higgs
- We are already probing $\tan\beta \sim 40$
- Combination of the results in all channels from CDF and D0 will result in a significant boost in sensitivity

"The hardest thing of all is to find a black cat in a dark room, especially if there is no cat."

- Confucius